

ON THE INDEX OF A NON-FREDHOLM MODEL OPERATOR

ALAN CAREY, FRITZ GESZTESY, GALINA LEVITINA, AND FEDOR SUKOCHEV

Dedicated to the memory of Leiba Rodman (1949–2015).

ABSTRACT. Let $\{A(t)\}_{t \in \mathbb{R}}$ be a path of self-adjoint Fredholm operators in a Hilbert space \mathcal{H} , joining endpoints A_{\pm} as $t \rightarrow \pm\infty$. Computing the index of the operator $\mathbf{D}_A = \partial/\partial t + A$ acting on $L^2(\mathbb{R}; \mathcal{H})$, where A denotes the multiplication operator $(Af)(t) = A(t)f(t)$ for $f \in L^2(\mathbb{R}; \mathcal{H})$, and its relation to spectral flow along this path, has a long history, but it is mostly focussed on the case where the operators $A(t)$ all have purely discrete spectrum.

Introducing the operators $\mathbf{H}_1 = \mathbf{D}_A^* \mathbf{D}_A$ and $\mathbf{H}_2 = \mathbf{D}_A \mathbf{D}_A^*$, we consider spectral shift functions, denoted by $\xi(\cdot; A_+, A_-)$ and $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ associated with the pairs (A_+, A_-) and $(\mathbf{H}_2, \mathbf{H}_1)$. Under the restrictive hypotheses that A_+ is a relatively trace class perturbation of A_- , a relationship between these spectral shift functions was proved in [14], for certain operators A_{\pm} with essential spectrum, extending a result of Pushnitski [22]. Moreover, assuming A_{\pm} to be Fredholm, the value $\xi(0; A_-, A_+)$ was shown to represent the spectral flow along the path $\{A(t)\}_{t \in \mathbb{R}}$ while that of $\xi(0_+; \mathbf{H}_1, \mathbf{H}_2)$ yields the Fredholm index of \mathbf{D}_A . The fact, proved in [14], that these values of the two spectral functions are equal, resolves the index = spectral flow question in this case. This relationship between spectral shift functions was generalized to non-Fredholm operators in [9] again under the relatively trace class perturbation hypothesis. In this situation it asserts that the Witten index of \mathbf{D}_A , denoted by $W_r(\mathbf{D}_A)$, a substitute for the Fredholm index in the absence of the Fredholm property of \mathbf{D}_A , is given by

$$W_r(\mathbf{D}_A) = \xi_L(0_+; \mathbf{H}_2, \mathbf{H}_1) = [\xi_L(0_+; A_+, A_-) + \xi_L(0_-; A_+, A_-)]/2.$$

Here one assumes that $\xi(\cdot; A_-, A_+)$ possesses a right and left Lebesgue point at 0 denoted by $\xi_L(0_{\pm}; A_+, A_-)$ (and similarly for $\xi_L(0_+; \mathbf{H}_2, \mathbf{H}_1)$).

When the path $\{A(t)\}_{t \in \mathbb{R}}$ consists of differential operators, the relatively trace class perturbation assumption is violated. The simplest assumption that applies (to differential operators in 1+1 dimensions) is to admit relatively Hilbert–Schmidt perturbations. This is not just an incremental improvement. In fact, the method we employ here to make this extension is of interest in any dimension. Moreover we consider A_{\pm} which are not necessarily Fredholm and we establish that the relationships between the two spectral shift functions found in **all** of the previous papers [9], [14], and [22], can be proved, even in the non-Fredholm case. The significance of our new methods is that, besides being simpler, they also allow a wide class of examples such as pseudodifferential operators in higher dimensions. Most importantly, we prove the above formula for the Witten index in the most general circumstances to date.

Date: September 8, 2015.

2010 Mathematics Subject Classification. Primary 47A53, 58J30; Secondary 47A10, 47A40.

Key words and phrases. Fredholm and Witten index, spectral shift function.

To appear in *Operators and Matrices*.

CONTENTS

1. Introduction	2
2. The Strategy Employed	4
3. The Basic Setup	6
4. Computing $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ In Terms Of $\xi(\cdot; A_+, A_-)$	17
5. The Witten Index	19
6. A $(1+1)$ -Dimensional Example	20
Appendix A. Some Facts On Spectral Shift Functions, Trace Formulas, and Modified Fredholm Determinants	23
References	28

1. INTRODUCTION

Typical Hamiltonians in quantum mechanical models have some essential spectrum and this is true more generally for differential operators on non-compact manifolds. Much less is known about Fredholm theory and spectral flow in these situations. Our objective in this paper is to investigate operators for which zero is in the essential spectrum and for which Fredholm theory is not applicable.

This paper is motivated by [9] where results on an index theory for certain non-Fredholm operators are described using the model operator formalism in [14]. The latter paper was inspired by [22] which, in turn, was motivated by [24], where the relationship between the Fredholm index and spectral flow for operators with discrete spectrum is studied. The model operators considered there provide prototypes for more complex situations. They arise in connection with investigations of the Maslov index, Morse theory, Floer homology, Sturm oscillation theory, etc.

The principal aim in [22] and [14] was to extend the discrete spectrum results of [24], relating the Fredholm index and spectral flow, to a relatively trace class perturbation theory approach, permitting essential spectra. However the relatively trace class assumption rules out standard differential operators such as Dirac-type operators and thus, in order to incorporate this important class of examples, we need a more general framework. In this paper we introduce a new approach that enables us to handle relatively Hilbert–Schmidt perturbations. This improvement incorporates $1+1$ dimensional differential operators while the methods introduced here are also applicable in more general situations. Moreover we prove the main result of [14] by a shorter and simpler method at the same time as generalizing it substantially so that it applies to non-Fredholm operators.

To introduce the model, let $\{A(t)\}_{t \in \mathbb{R}}$ be a family of self-adjoint operators in the complex, separable Hilbert space \mathcal{H} , subject to the assumption that self-adjoint limiting operators

$$A_+ = \lim_{t \rightarrow +\infty} A(t), \quad A_- = \lim_{t \rightarrow -\infty} A(t) \quad (1.1)$$

exist in \mathcal{H} in the norm resolvent sense. We denote by \mathbf{A} the operator in $L^2(\mathbb{R}; \mathcal{H})$ defined by

$$\begin{aligned} (\mathbf{A}f)(t) &= A(t)f(t) \text{ for a.e. } t \in \mathbb{R}, \\ f \in \text{dom}(\mathbf{A}) &= \left\{ g \in L^2(\mathbb{R}; \mathcal{H}) \left| \begin{aligned} &g(t) \in \text{dom}(A(t)) \text{ for a.e. } t \in \mathbb{R}; \\ &t \mapsto A(t)g(t) \text{ is (weakly) measurable; } \int_{\mathbb{R}} dt \|A(t)g(t)\|_{\mathcal{H}}^2 < \infty \end{aligned} \right. \right\}. \end{aligned} \quad (1.2)$$

Next, we introduce the model operator

$$\mathbf{D}_A = \frac{d}{dt} + A, \quad \text{dom}(\mathbf{D}_A) = \text{dom}(d/dt) \cap \text{dom}(A_-), \quad (1.3)$$

and the associated nonnegative, self-adjoint operators

$$\mathbf{H}_1 = \mathbf{D}_A^* \mathbf{D}_A, \quad \mathbf{H}_2 = \mathbf{D}_A \mathbf{D}_A^*, \quad (1.4)$$

in $L^2(\mathbb{R}; \mathcal{H})$. (Here A_- in $L^2(\mathbb{R}; \mathcal{H})$ represents the self-adjoint constant fiber operator defined according to (1.2), with $A(t)$ replaced by the asymptote A_-).

Assuming that A_- and A_+ are boundedly invertible, we also recall (cf. [14]) that \mathbf{D}_A is a Fredholm operator in $L^2(\mathbb{R}; \mathcal{H})$. Under the additional relatively trace class assumption, that is, $(A_+ - A_-)(A_-^2 + I_{\mathcal{H}})^{-1/2}$ is trace class, it is shown in [14] (and earlier in [22] under a more stringent set of hypotheses on the family $A(\cdot)$), that the Fredholm index of \mathbf{D}_A may then be computed as follows,

$$\text{index}(\mathbf{D}_A) = \xi(0_+; \mathbf{H}_2, \mathbf{H}_1) = \xi(0; A_+, A_-). \quad (1.5)$$

Here $\xi(\cdot; S_2, S_1)$ denotes the spectral shift function for the pair of self-adjoint operators (S_2, S_1) . Whenever S_j , $j = 1, 2$, are bounded from below, we adhere to the normalization

$$\xi(\lambda; S_2, S_1) = 0 \text{ for } \lambda < \inf(\sigma(S_1) \cup \sigma(S_2)), \quad (1.6)$$

in particular, $\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) = 0$, $\lambda < 0$.

The new direction developed in [9] focuses on the model operator \mathbf{D}_A in $L^2(\mathbb{R}; \mathcal{H})$ whenever the latter ceases to be Fredholm. First, we recall the definition of the Witten index as studied in [4] and [18]:

$$W_r(\mathbf{D}_A) = \lim_{\lambda \uparrow 0} (-\lambda) \text{tr}_{L^2(\mathbb{R}; \mathcal{H})} ((\mathbf{H}_1 - \lambda \mathbf{I})^{-1} - (\mathbf{H}_2 - \lambda \mathbf{I})^{-1}), \quad (1.7)$$

whenever the limit exists (where $\text{tr}_{\mathcal{K}}(\cdot)$ abbreviates the trace in the Hilbert space \mathcal{K}). Here, the subscript “r” refers to the resolvent regularization used; other regularizations, for instance, semigroup (heat kernel) based ones, are possible (cf., [9]).

If \mathbf{D}_A is Fredholm (and of course the necessary trace class conditions in (1.7) are satisfied), one has consistency with the Fredholm index of \mathbf{D}_A . In addition, under appropriate spectral assumptions the following connection between Fredholm, respectively, Witten indices and the underlying spectral shift function applies

$$W_r(\mathbf{D}_A) = \xi(0_+; \mathbf{H}_2, \mathbf{H}_1). \quad (1.8)$$

Most importantly, $W_r(\mathbf{D}_A)$ exhibits invariance properties under additive, relatively trace class perturbations (apart from some additional technical hypotheses). This is sometimes dubbed topological invariance of the Witten index in the pertinent literature (see, e.g., [4], [5], [14], [18], and the references therein).

Originally, index regularizations such as (1.7) were studied in the context of supersymmetric quantum mechanics in the physics literature in the 1970’s and 1980’s, see, [9] for details.

The results of [9] for the specific model operator \mathbf{D}_A in $L^2(\mathbb{R}; \mathcal{H})$ are as follows. Assuming that 0 is a right and a left Lebesgue point of $\xi(\cdot; A_+, A_-)$ (denoted by $\xi_L(0_+; A_+, A_-)$ and $\xi_L(0_-; A_+, A_-)$, respectively), then it is also a right Lebesgue point of $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ (denoted by $\xi_L(0_+; \mathbf{H}_2, \mathbf{H}_1)$). Under this right/left Lebesgue point assumption on $\xi(\cdot; A_+, A_-)$, the analog of (1.3) and (1.4) in the general case

where D_A ceases to be Fredholm, the principal new result of [9], then reads as follows,

$$W_r(D_A) = \xi_L(0_+; H_2, H_1) = [\xi_L(0_+; A_+, A_-) + \xi_L(0_-; A_+, A_-)]/2. \quad (1.9)$$

Recalling our earlier mention of Dirac operators, let A_- be the flat space Dirac operator on spinor valued functions on \mathbb{R}^n and perturb it by a multiplication operator by a function $\phi : \mathbb{R}^n \rightarrow \mathbb{R}$ (matrix-valued functions also provide examples) to define the operator $A_+ = A_- + \phi$. It follows from the discussion in Remark (c) of [25, Chapter 4] that, for ϕ of sufficiently rapid decay at $\pm\infty$, we can show that $(A_+ - A_-)(A_-^2 + I_{\mathcal{H}})^{-s/2}$ is trace class for $s > n$ but for no lesser value of s . Thus, even in one dimension the relatively trace class perturbation assumption is violated for geometric examples based on Dirac-type operators.

Our main objective in this paper is to extend the results described above to situations in which the relatively trace class perturbation assumption no longer holds, replacing it by Hypothesis 3.10 below. The new technique described here is an approximation argument that amounts, for differential operators, to using pseudodifferential approximating operators. Then, for the approximants, the relatively trace class perturbation condition is restored and the results of our earlier papers on spectral shift functions ([8], [9]) are available for use. So we obtain strong information on the spectral shift functions for the approximants. Then we find that they may be shown to converge to the spectral shift functions for the original operators, as the pseudodifferential approximating operators converge in an appropriate (strong resp., norm resolvent) sense. This leads to one of our main results, (1.9) under the most general Hypothesis (3.10) to date, which generalizes the main theorem of [14] and the subsequent [9].

While the strategy of this paper applies in higher dimensions (where the relatively Hilbert–Schmidt condition is replaced by a relatively Schatten class condition) further additional ideas are needed to make our approximation scheme work there. Examples in [10] illustrate some of the issues. This matter is a part of ongoing investigations. We remark that the topological meaning of the Witten index is explored in [18] and [11] while its geometric significance is still under investigation.

2. THE STRATEGY EMPLOYED

In this section we briefly outline the principal new strategy employed in this paper that permits us to circumvent the relative trace class Hypothesis 2.1 used in [14] and [9]. Throughout this section we assume that Hypothesis 3.10 is valid.

I. Consider the family of self-adjoint operators $A(t)$, $t \in \mathbb{R}$, with asymptotes A_{\pm} as well as $B(t)$, $t \in \mathbb{R}$, and B_{\pm} , such that

$$A(t) = A_- + B(t), \quad t \in \mathbb{R}. \quad (2.1)$$

Introduce $D_A = \frac{d}{dt} + A$ in $L^2(\mathbb{R}; \mathcal{H})$ as in (1.3), and define A_- in $L^2(\mathbb{R}; \mathcal{H})$ as the self-adjoint (constant fiber) operator defined in (3.14). Next, introduce the operators H_j , $j = 1, 2$, in $L^2(\mathbb{R}; \mathcal{H})$ by $H_1 = D_A^* D_A$, $H_2 = D_A D_A^*$ as in (1.4). In addition, assuming continuous differentiability of the family $B(\cdot)$ and introducing B and B' in terms of the bounded operator families $B(t)$, $B'(t)$, $t \in \mathbb{R}$, in analogy to (2.16), one can decompose H_j , $j = 1, 2$, as follows:

$$H_j = H_0 + B A_- + A_- B + B^2 + (-1)^j B', \quad j = 1, 2, \quad (2.2)$$

with

$$\mathbf{A} = \mathbf{A}_- + \mathbf{B}, \quad \text{dom}(\mathbf{A}) = \text{dom}(\mathbf{A}_-). \quad (2.3)$$

II. Use the approximation

$$A_n(t) = A_- + \chi_n(A_-)B(t)\chi_n(A_-), \quad n \in \mathbb{N}, \quad t \in \mathbb{R}, \quad (2.4)$$

with

$$\chi_n(\nu) = \frac{n}{(\nu^2 + n^2)^{1/2}}, \quad \nu \in \mathbb{R}, \quad n \in \mathbb{N}, \quad (2.5)$$

such that $\text{s-lim}_{n \rightarrow \infty} \chi_n(A_-) = I$.

Given our assumptions on $B(\cdot)$ (cf. Hypothesis 3.10) one infers that

$$A_{-,n} = A_-, \quad A_{+,n} = A_- + \chi_n(A_-)B_+\chi_n(A_-), \quad n \in \mathbb{N}, \quad (2.6)$$

$$A_{+,n} - A_- = \chi_n(A_-)B_+\chi_n(A_-) \in \mathcal{B}_1(\mathcal{H}), \quad n \in \mathbb{N}, \quad (2.7)$$

$$A'_n(t) = B'_n(t) = \chi_n(A_-)B'(t)\chi_n(A_-) \in \mathcal{B}_1(\mathcal{H}), \quad n \in \mathbb{N}, \quad t \in \mathbb{R}, \quad (2.8)$$

$$\int_{\mathbb{R}} dt \|A'_n(t)\|_{\mathcal{B}_1(\mathcal{H})} < \infty. \quad (2.9)$$

(We denote by $\mathcal{B}_p(\mathcal{H})$ the standard ℓ^p -based trace ideals, $p \geq 1$.) Thus, one also obtains,

$$\begin{aligned} \mathbf{H}_{j,n} &= \mathbf{H}_0 + \mathbf{B}_n \mathbf{A}_- + \mathbf{A}_- \mathbf{B}_n + \mathbf{B}_n^2 + (-1)^j \mathbf{B}'_n, \quad \text{dom}(\mathbf{H}_{j,n}) = \text{dom}(\mathbf{H}_0), \\ &n \in \mathbb{N}, \quad j = 1, 2, \end{aligned} \quad (2.10)$$

with

$$\mathbf{B}_n = \chi_n(\mathbf{A}_-)\mathbf{B}\chi_n(\mathbf{A}_-), \quad \mathbf{B}'_n = \chi_n(\mathbf{A}_-)\mathbf{B}'\chi_n(\mathbf{A}_-), \quad n \in \mathbb{N}. \quad (2.11)$$

III. As a consequence of step **II**, the spectral shift functions $\xi(\cdot; A_{+,n}, A_-)$ and $\xi(\cdot; \mathbf{H}_{2,n}, \mathbf{H}_{1,n})$, $n \in \mathbb{N}$, exist and are uniquely determined by

$$\xi(\cdot; A_{+,n}, A_-) \in L^1(\mathbb{R}; d\nu), \quad \xi(\lambda; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) = 0, \quad \lambda < 0, \quad n \in \mathbb{N}. \quad (2.12)$$

Moreover, employing [8] or [22], one obtains the approximate trace formula,

$$\int_{[0,\infty)} \frac{\xi(\lambda; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) d\lambda}{(\lambda - z)^2} = \frac{1}{2} \int_{\mathbb{R}} \frac{\xi(\nu; A_{+,n}, A_-) d\nu}{(\nu^2 - z)^{3/2}}, \quad n \in \mathbb{N}, \quad z \in \mathbb{C} \setminus [0, \infty). \quad (2.13)$$

IV. The main result. Now we take the limits $n \rightarrow \infty$ in (2.13). We use the trace norm convergence result in Theorem 3.13 in combination with some Fredholm determinant facts to control the limit $n \rightarrow \infty$ of the left-hand and right-hand side of (2.13) to arrive at

$$\int_{[0,\infty)} \frac{\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) d\lambda}{(\lambda - z)^2} = \frac{1}{2} \int_{\mathbb{R}} \frac{\xi(\nu; A_+, A_-) d\nu}{(\nu^2 - z)^{3/2}}, \quad z \in \mathbb{C} \setminus [0, \infty). \quad (2.14)$$

Relation (2.14) combined with a Stieltjes inversion argument then implies the main formula of the paper (a Pushnitski-type relation between spectral shift functions):

$$\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) = \frac{1}{\pi} \int_{-\lambda^{1/2}}^{\lambda^{1/2}} \frac{\xi(\nu; A_+, A_-) d\nu}{(\lambda - \nu^2)^{1/2}} \quad \text{for a.e. } \lambda > 0. \quad (2.15)$$

As a result, one of the principal theorems proven in this paper then proves the equalities in (1.9). In particular, Equations (2.14), (2.15), and (1.9) represent the analog of the principal results in [14] and [9] for the model operator \mathbf{D}_A , but now under considerably more general hypotheses.

Notation. We briefly summarize some of the notation used in this paper: Let \mathcal{H} be a separable complex Hilbert space, $(\cdot, \cdot)_{\mathcal{H}}$ the scalar product in \mathcal{H} (linear in the second argument), and $I_{\mathcal{H}}$ the identity operator in \mathcal{H} .

Next, if T is a linear operator mapping (a subspace of) a Hilbert space into another, then $\text{dom}(T)$ and $\ker(T)$ denote the domain and kernel (i.e., null space) of T . The closure of a closable operator S is denoted by \overline{S} . The spectrum, essential spectrum, discrete spectrum, point spectrum, and resolvent set of a closed linear operator in a Hilbert space will be denoted by $\sigma(\cdot)$, $\sigma_{\text{ess}}(\cdot)$, $\sigma_d(\cdot)$, $\sigma_p(\cdot)$, and $\rho(\cdot)$, respectively.

The convergence of bounded operators in the strong operator topology (i.e., pointwise limits) will be denoted by s-lim, similarly, norm limits of bounded operators are denoted by n-lim.

The Banach spaces of bounded and compact linear operators on a separable complex Hilbert space \mathcal{H} are denoted by $\mathcal{B}(\mathcal{H})$ and $\mathcal{B}_{\infty}(\mathcal{H})$, respectively; the corresponding ℓ^p -based trace ideals will be denoted by $\mathcal{B}_p(\mathcal{H})$, $p > 0$. Moreover, $\det_{\mathcal{H}}(I_{\mathcal{H}} - A)$, and $\text{tr}_{\mathcal{H}}(A)$ denote the standard Fredholm determinant and the corresponding trace of a trace class operator $A \in \mathcal{B}_1(\mathcal{H})$.

Linear operators in the Hilbert space $L^2(\mathbb{R}; dt; \mathcal{H})$, in short, $L^2(\mathbb{R}; \mathcal{H})$, will be denoted by calligraphic boldface symbols of the type \mathbf{T} , to distinguish them from operators T in \mathcal{H} . In particular, operators denoted by \mathbf{T} in the Hilbert space $L^2(\mathbb{R}; \mathcal{H})$ typically represent operators associated with a family of operators $\{T(t)\}_{t \in \mathbb{R}}$ in \mathcal{H} , defined by

$$\begin{aligned} (\mathbf{T}f)(t) &= T(t)f(t) \text{ for a.e. } t \in \mathbb{R}, \\ f \in \text{dom}(\mathbf{T}) &= \left\{ g \in L^2(\mathbb{R}; \mathcal{H}) \left| g(t) \in \text{dom}(T(t)) \text{ for a.e. } t \in \mathbb{R}; \right. \right. \\ &\quad \left. \left. t \mapsto T(t)g(t) \text{ is (weakly) measurable; } \int_{\mathbb{R}} dt \|T(t)g(t)\|_{\mathcal{H}}^2 < \infty \right\}. \end{aligned} \quad (2.16)$$

In the special case, where $\{T(t)\}$ is a family of bounded operators on \mathcal{H} with $\sup_{t \in \mathbb{R}} \|T(t)\|_{\mathcal{B}(\mathcal{H})} < \infty$, the associated operator \mathbf{T} is a bounded operator on $L^2(\mathbb{R}; \mathcal{H})$ with $\|\mathbf{T}\|_{\mathcal{B}(L^2(\mathbb{R}; \mathcal{H}))} = \sup_{t \in \mathbb{R}} \|T(t)\|_{\mathcal{B}(\mathcal{H})}$.

For brevity we will abbreviate $I := I_{\mathcal{H}}$ and $\mathbf{I} := I_{L^2(\mathbb{R}; \mathcal{H})}$.

3. THE BASIC SETUP

In this section we set the stage for an extension of [14] that circumvents the apparently fundamental relative trace class condition in Hypothesis 2.1 in [14] and [9].

We start by collecting our principal assumptions:

Hypothesis 3.1. *Suppose \mathcal{H} is a complex, separable Hilbert space.*

- (i) *Assume A_- is self-adjoint on $\text{dom}(A_-) \subseteq \mathcal{H}$.*
- (ii) *Suppose we have a family of bounded operators $\{B(t)\}_{t \in \mathbb{R}} \subset \mathcal{B}(\mathcal{H})$, continuously differentiable in norm on \mathbb{R} such that*

$$\|B'(\cdot)\|_{\mathcal{B}(\mathcal{H})} \in L^1(\mathbb{R}; dt). \quad (3.1)$$

Given Hypothesis 3.1, we introduce the family of self-adjoint operators $A(t)$, $t \in \mathbb{R}$, in \mathcal{H} , by

$$A(t) = A_- + B(t), \quad \text{dom}(A(t)) = \text{dom}(A_-), \quad t \in \mathbb{R}. \quad (3.2)$$

Moreover, writing

$$B(t) = B(t_0) + \int_{t_0}^t ds B'(s), \quad t, t_0 \in \mathbb{R}, \quad (3.3)$$

one infers that the self-adjoint asymptotes

$$\text{n-lim}_{t \rightarrow \pm\infty} B(t) := B_{\pm} \in \mathcal{B}(\mathcal{H}) \quad (3.4)$$

exist. In particular, we will make the choice

$$B_- = 0 \quad (3.5)$$

in the following and also introduce the asymptote,

$$A_+ = A_- + B_+, \quad \text{dom}(A_+) = \text{dom}(A_-). \quad (3.6)$$

Assumption (3.1) also yields,

$$\sup_{t \in \mathbb{R}} \|B(t)\|_{\mathcal{B}(\mathcal{H})} \leq \int_{\mathbb{R}} dt \|B'(t)\|_{\mathcal{B}(\mathcal{H})} < \infty. \quad (3.7)$$

A simple application of the resolvent identity yields (with $t \in \mathbb{R}$, $z \in \mathbb{C} \setminus \mathbb{R}$)

$$(A(t) - zI)^{-1} = (A_{\pm} - zI)^{-1} - (A(t) - zI)^{-1}[B(t) - B_{\pm}](A_{\pm} - zI)^{-1}, \quad (3.8)$$

$$\|(A(t) - zI)^{-1} - (A_{\pm} - zI)^{-1}\|_{\mathcal{B}(\mathcal{H})} \leq |\text{Im}(z)|^{-2} \|B(t) - B_{\pm}\|_{\mathcal{B}(\mathcal{H})}, \quad (3.9)$$

and hence proves that

$$\text{n-lim}_{t \rightarrow \pm\infty} (A(t) - zI)^{-1} = (A_{\pm} - zI)^{-1}, \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (3.10)$$

At this point we need to introduce additional hypotheses to those in Hypothesis 3.1. These additional requirements can be accommodated for differential operators in low dimensions. We know that a weakening of the next hypothesis is needed for higher-dimensional geometric examples such as those provided by Dirac-type operators.

Hypothesis 3.2. *In addition to Hypothesis 3.1, assume the following conditions on B_+ , $B(t)$, $t \in \mathbb{R}$:*

(i) *Suppose that $|B_+|^{1/2}(A_- - z_0 I)^{-1} \in \mathcal{B}_2(\mathcal{H})$ for some (and hence for all) $z_0 \in \rho(A_-)$.*

(ii) *Assume that $\sup_{t \in \mathbb{R}} \|B'(t)\|_{\mathcal{B}(\mathcal{H})} < \infty$.*

Remark 3.3. For Dirac-type operators, the Hilbert–Schmidt condition in Hypothesis 3.2 is directly tied to the fact that we are considering differences of resolvents in (3.31). We know from [10] that we need to consider differences of higher powers of resolvents in higher-dimensional examples whose treatment is deferred to future investigations. \diamond

Assuming Hypothesis 3.2 in the following, the resolvent identity

$$(A_+ - zI)^{-1} = (A_- - zI)^{-1} - (A_+ - zI)^{-1}B_+(A_- - zI)^{-1}, \quad z \in \mathbb{C} \setminus \mathbb{R}, \quad (3.11)$$

combined with Hypothesis 3.2 (i) yields

$$\begin{aligned} (A_+ - zI)^{-1} - (A_- - zI)^{-1} &= [(A_- - \bar{z}I)(A_+ - \bar{z}I)^{-1}]^* [|B_+|^{1/2}(A_- - \bar{z}I)^{-1}]^* \\ &\quad \times \text{sgn}(B_+) [|B_+|^{1/2}(A_- - zI)^{-1}], \quad z \in \mathbb{C} \setminus \mathbb{R}, \end{aligned} \quad (3.12)$$

and hence

$$[(A_+ - zI)^{-1} - (A_- - zI)^{-1}] \in \mathcal{B}_1(\mathcal{H}), \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (3.13)$$

Next, we turn to the pair $(\mathbf{H}_2, \mathbf{H}_1)$, assuming Hypothesis 3.2: First, we recall that $\mathbf{A}, \mathbf{B}, \mathbf{A}' = \mathbf{B}'$, are defined in terms of the families $A(t)$, $B(t)$, and $B'(t)$, $t \in \mathbb{R}$, as in (2.16). In addition, \mathbf{A}_- in $L^2(\mathbb{R}; \mathcal{H})$ represents the self-adjoint (constant fiber) operator defined by

$$\begin{aligned} (\mathbf{A}_- f)(t) &= A_- f(t) \text{ for a.e. } t \in \mathbb{R}, \\ f \in \text{dom}(\mathbf{A}_-) &= \left\{ g \in L^2(\mathbb{R}; \mathcal{H}) \left| g(t) \in \text{dom}(A_-) \text{ for a.e. } t \in \mathbb{R}, \right. \right. \\ &\quad \left. \left. t \mapsto A_- g(t) \text{ is (weakly) measurable, } \int_{\mathbb{R}} dt \|A_- g(t)\|_{\mathcal{H}}^2 < \infty \right\}. \end{aligned} \quad (3.14)$$

Now we introduce the operator $\mathbf{D}_\mathbf{A}$ in $L^2(\mathbb{R}; \mathcal{H})$ by

$$\mathbf{D}_\mathbf{A} = \frac{d}{dt} + \mathbf{A}, \quad \text{dom}(\mathbf{D}_\mathbf{A}) = W^{1,2}(\mathbb{R}; \mathcal{H}) \cap \text{dom}(\mathbf{A}_-), \quad (3.15)$$

where we used that (3.7) implies

$$\|\mathbf{B}\|_{\mathcal{B}(L^2(\mathbb{R}; \mathcal{H}))} = \sup_{t \in \mathbb{R}} \|B(t)\|_{\mathcal{B}(\mathcal{H})} < \infty \quad (3.16)$$

and hence

$$\mathbf{A} = \mathbf{A}_- + \mathbf{B}, \quad \text{dom}(\mathbf{A}) = \text{dom}(\mathbf{A}_-). \quad (3.17)$$

Here the operator d/dt in $L^2(\mathbb{R}; \mathcal{H})$ is defined by

$$\begin{aligned} \left(\frac{d}{dt} f \right)(t) &= f'(t) \text{ for a.e. } t \in \mathbb{R}, \\ f \in \text{dom}(d/dt) &= \{ g \in L^2(\mathbb{R}; \mathcal{H}) \mid g \in AC_{\text{loc}}(\mathbb{R}; \mathcal{H}), g' \in L^2(\mathbb{R}; \mathcal{H}) \} \\ &= W^{1,2}(\mathbb{R}; \mathcal{H}). \end{aligned} \quad (3.18)$$

Clearly (cf. [14, Lemma 4.4], which also holds for our more general setting), $\mathbf{D}_\mathbf{A}$ is densely defined and closed in $L^2(\mathbb{R}; \mathcal{H})$ and the adjoint operator $\mathbf{D}_\mathbf{A}^*$ of $\mathbf{D}_\mathbf{A}$ is given by

$$\mathbf{D}_\mathbf{A}^* = -\frac{d}{dt} + \mathbf{A}, \quad \text{dom}(\mathbf{D}_\mathbf{A}^*) = W^{1,2}(\mathbb{R}; \mathcal{H}) \cap \text{dom}(\mathbf{A}_-). \quad (3.19)$$

This enables one to introduce the nonnegative, self-adjoint operators \mathbf{H}_j , $j = 1, 2$, in $L^2(\mathbb{R}; \mathcal{H})$ by

$$\mathbf{H}_1 = \mathbf{D}_\mathbf{A}^* \mathbf{D}_\mathbf{A}, \quad \mathbf{H}_2 = \mathbf{D}_\mathbf{A} \mathbf{D}_\mathbf{A}^*. \quad (3.20)$$

In order to effectively describe the domains of \mathbf{H}_j , $j = 1, 2$, we need to decompose the latter as follows: First, we introduce \mathbf{B}' in terms of the bounded operator families $B'(t)$, $t \in \mathbb{R}$, in analogy to (2.16), and observe that Hypothesis 3.2 (ii) implies that

$$\|\mathbf{B}'\|_{\mathcal{B}(L^2(\mathbb{R}; \mathcal{H}))} = \sup_{t \in \mathbb{R}} \|B'(t)\|_{\mathcal{B}(\mathcal{H})} < \infty. \quad (3.21)$$

Next, we strengthen our hypothesis on $B(t)$, $t \in \mathbb{R}$, as follows. Introduce \mathbf{H}_0 in $L^2(\mathbb{R}; \mathcal{H})$ by

$$\mathbf{H}_0 = -\frac{d^2}{dt^2} + \mathbf{A}_-^2, \quad \text{dom}(\mathbf{H}_0) = W^{2,2}(\mathbb{R}; \mathcal{H}) \cap \text{dom}(\mathbf{A}_-^2). \quad (3.22)$$

Then \mathbf{H}_0 is self-adjoint by Theorem VIII.33 of [23].

Again we need to make some additional hypotheses motivated by differential operator examples.

Hypothesis 3.4. *In addition to Hypotheses 3.2, assume the following conditions on $B(t)$, $t \in \mathbb{R}$:*

(i) *Suppose that $\mathbf{A}_-\mathbf{B}$ is bounded with respect to \mathbf{H}_0 with bound strictly less than one, that is, there exists $a \in (0, 1)$ and $b \in (0, \infty)$ such that*

$$\|\mathbf{A}_-\mathbf{B}f\|_{L^2(\mathbb{R}; \mathcal{H})} \leq a\|\mathbf{H}_0 f\|_{L^2(\mathbb{R}; \mathcal{H})} + b\|f\|_{L^2(\mathbb{R}; \mathcal{H})}, \quad f \in \text{dom}(\mathbf{H}_0). \quad (3.23)$$

(ii) *Assume that for some (and hence for all) $z_0 \in \mathbb{C} \setminus \mathbb{R}$,*

$$|\mathbf{B}'|^{1/2}(\mathbf{H}_0 - z_0 \mathbf{I})^{-1} \in \mathcal{B}_2(L^2(\mathbb{R}; \mathcal{H})) \quad (3.24)$$

Remark 3.5. While it is clear that $\mathbf{B}\mathbf{A}_-$ is infinitesimally bounded with respect to \mathbf{H}_0 , to prove this it suffices to note that

$$\begin{aligned} \|\mathbf{B}\mathbf{A}_-(\mathbf{H}_0 - z\mathbf{I})^{-1}\|_{\mathcal{B}(L^2(\mathbb{R}; \mathcal{H}))} &\leq |\text{Im}(z)|^{-1/2}\|\mathbf{B}\|_{\mathcal{B}(L^2(\mathbb{R}; \mathcal{H}))} \\ &\times \|\mathbf{A}_-(\mathbf{H}_0 - z\mathbf{I})^{-1/2}\|_{\mathcal{B}(L^2(\mathbb{R}; \mathcal{H}))}, \quad z \in \mathbb{C} \setminus [0, \infty), \end{aligned} \quad (3.25)$$

it is not obvious that $\mathbf{A}_-\mathbf{B}$ is bounded with respect to \mathbf{H}_0 . For later purpose we note that (3.23) implies the existence of $a' \in (a, 1)$ such that

$$\|\mathbf{A}_-\mathbf{B}(\mathbf{H}_0 - z\mathbf{I})^{-1}\|_{\mathcal{B}(L^2(\mathbb{R}; \mathcal{H}))} \leq a' < 1 \quad (3.26)$$

for $0 < |\text{Im}(z)|$ sufficiently large. \diamond

Assuming Hypothesis 3.4 in the following, (3.23) combined with (3.25) imply that the operator $\mathbf{B}\mathbf{A}_- + \mathbf{A}_-\mathbf{B}$ is \mathbf{H}_0 -bounded with bound less than one, and therefore, by [21, Theorem VI.4.3] the following decomposition of the operators \mathbf{H}_j , $j = 1, 2$ holds

$$\begin{aligned} \mathbf{H}_j &= \frac{d^2}{dt^2} + \mathbf{A}^2 + (-1)^j \mathbf{A}' \\ &= \mathbf{H}_0 + \mathbf{B}\mathbf{A}_- + \mathbf{A}_-\mathbf{B} + \mathbf{B}^2 + (-1)^j \mathbf{B}', \\ \text{dom}(\mathbf{H}_j) &= \text{dom}(\mathbf{H}_0), \quad j = 1, 2. \end{aligned} \quad (3.27)$$

In addition,

$$\begin{aligned} (\mathbf{H}_2 - z\mathbf{I})^{-1} - (\mathbf{H}_1 - z\mathbf{I})^{-1} &= -(\mathbf{H}_1 - z\mathbf{I})^{-1}[2\mathbf{B}'](\mathbf{H}_2 - z\mathbf{I})^{-1}, \\ &= -2[|\mathbf{B}'|^{1/2}(\mathbf{H}_1 - \bar{z}\mathbf{I})^{-1}]^* \text{sgn}(\mathbf{B}')|\mathbf{B}'|^{1/2}(\mathbf{H}_2 - z\mathbf{I})^{-1}, \\ &\quad z \in \rho(\mathbf{H}_1) \cap \rho(\mathbf{H}_2), \end{aligned} \quad (3.28)$$

and using that

$$\overline{(\mathbf{H}_1 - z\mathbf{I})^{-1}(\mathbf{H}_0 - z\mathbf{I})} = [(\mathbf{H}_0 - \bar{z}\mathbf{I})(\mathbf{H}_1 - \bar{z}\mathbf{I})^{-1}]^* \in \mathcal{B}(L^2(\mathbb{R}; \mathcal{H})), \quad (3.29)$$

$$(\mathbf{H}_0 - z\mathbf{I})(\mathbf{H}_2 - z\mathbf{I})^{-1} \in \mathcal{B}(L^2(\mathbb{R}; \mathcal{H})), \quad z \in \rho(\mathbf{H}_1) \cap \rho(\mathbf{H}_2), \quad (3.30)$$

and assumption of Hypothesis 3.4 (ii) one concludes that

$$[(\mathbf{H}_2 - z\mathbf{I})^{-1} - (\mathbf{H}_1 - z\mathbf{I})^{-1}] \in \mathcal{B}_1(L^2(\mathbb{R}; \mathcal{H})), \quad z \in \rho(\mathbf{H}_1) \cap \rho(\mathbf{H}_2). \quad (3.31)$$

The fact (3.31) implies that the spectral shift function $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ for the pair $(\mathbf{H}_2, \mathbf{H}_1)$ is well-defined, satisfies

$$\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1) \in L^1(\mathbb{R}; (\lambda^2 + 1)^{-1} d\lambda), \quad (3.32)$$

and since $\mathbf{H}_j \geq 0$, $j = 1, 2$, one uniquely introduces $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ by requiring that

$$\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) = 0, \quad \lambda < 0, \quad (3.33)$$

implying the Krein–Lifshitz trace formula,

$$\mathrm{tr}_{L^2(\mathbb{R}; \mathcal{H})} ((\mathbf{H}_2 - z\mathbf{I})^{-1} - (\mathbf{H}_1 - z\mathbf{I})^{-1}) = - \int_{[0, \infty)} \frac{\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) d\lambda}{(\lambda - z)^2}, \quad (3.34)$$

$$z \in \mathbb{C} \setminus [0, \infty).$$

Next, we deviate from the approximation procedure originally employed in [14] and [22]. We now introduce

$$\chi_n(\nu) = \frac{n}{(\nu^2 + n^2)^{1/2}}, \quad \nu \in \mathbb{R}, \quad n \in \mathbb{N}, \quad (3.35)$$

and hence obtain

$$\mathrm{s}\text{-}\lim_{n \rightarrow \infty} \chi_n(A_-) = I, \quad (3.36)$$

by an elementary application of the spectral theorem for A_- . The precise form of χ_n is of course immaterial, we just need property (3.36) and property (3.67) below.

We introduce

$$\begin{aligned} A_n(t) &= A_- + \chi_n(A_-)B(t)\chi_n(A_-) = A_- + B_n(t), \\ \mathrm{dom}(A_n(t)) &= \mathrm{dom}(A_-), \quad n \in \mathbb{N}, \quad t \in \mathbb{R}, \end{aligned} \quad (3.37)$$

where

$$B_n(t) = \chi_n(A_-)B(t)\chi_n(A_-), \quad n \in \mathbb{N}, \quad t \in \mathbb{R}. \quad (3.38)$$

In addition, we introduce

$$A_{-,n} = A_-, \quad \mathrm{dom}(A_{-,n}) = \mathrm{dom}(A_-), \quad n \in \mathbb{N}, \quad (3.39)$$

and conclude that

$$A_{+,n} = A_- + \chi_n(A_-)B_+\chi_n(A_-), \quad \mathrm{dom}(A_{+,n}) = \mathrm{dom}(A_-), \quad n \in \mathbb{N}, \quad (3.40)$$

$$A_{+,n} - A_- = \chi_n(A_-)B_+\chi_n(A_-) \in \mathcal{B}_1(\mathcal{H}), \quad n \in \mathbb{N}, \quad (3.41)$$

$$A'_n(t) = B'_n(t) = \chi_n(A_-)B'(t)\chi_n(A_-) \in \mathcal{B}_1(\mathcal{H}), \quad n \in \mathbb{N}, \quad t \in \mathbb{R}. \quad (3.42)$$

As a consequence of (3.41), the spectral shift functions $\xi(\cdot; A_{+,n}, A_-)$, $n \in \mathbb{N}$, exist and are uniquely determined by

$$\xi(\cdot; A_{+,n}, A_-) \in L^1(\mathbb{R}; d\nu), \quad n \in \mathbb{N}. \quad (3.43)$$

As a preparation to study various limits in Schatten–von Neumann ideals, we now recall the following standard convergence property for trace ideals:

Lemma 3.6. *Let $p \in [1, \infty)$ and assume that $R, R_n, T, T_n \in \mathcal{B}(\mathcal{H})$, $n \in \mathbb{N}$, satisfy $\mathrm{s}\text{-}\lim_{n \rightarrow \infty} R_n = R$ and $\mathrm{s}\text{-}\lim_{n \rightarrow \infty} T_n = T$ and that $S, S_n \in \mathcal{B}_p(\mathcal{H})$, $n \in \mathbb{N}$, satisfy $\lim_{n \rightarrow \infty} \|S_n - S\|_{\mathcal{B}_p(\mathcal{H})} = 0$. Then $\lim_{n \rightarrow \infty} \|R_n S_n T_n^* - R S T^*\|_{\mathcal{B}_p(\mathcal{H})} = 0$.*

As a first of several convergence results we state the following useful fact.

Lemma 3.7. *Assume Hypotheses 3.1 and 3.2(i). Then,*

$$\lim_{n \rightarrow \infty} \|(A_{+,n} - zI)^{-1} - (A_+ - zI)^{-1}\|_{\mathcal{B}_1(\mathcal{H})} = 0, \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (3.44)$$

Proof. One writes

$$\begin{aligned} (A_{+,n} - zI)^{-1} - (A_+ - zI)^{-1} &= [(A_{+,n} - zI)^{-1} - (A_- - zI)^{-1}] \\ &\quad - [(A_+ - zI)^{-1} - (A_- - zI)^{-1}] \\ &= -\chi_n(A_-)(A_- - zI)^{-1}B_+(A_- - zI)^{-1}\chi_n(A_-)[(A_- - zI)(A_{+,n} - zI)^{-1}] \\ &\quad + (A_- - zI)^{-1}B_+(A_- - zI)^{-1}[(A_- - zI)(A_+ - zI)^{-1}] \end{aligned}$$

$$\begin{aligned}
&= -\chi_n(A_-)(A_- - zI)^{-1}B_+(A_- - zI)^{-1}\chi_n(A_-) \\
&\quad \times [I - \chi_n(A_-)B_+\chi_n(A_-)(A_{+,n} - zI)^{-1}] \\
&\quad + (A_- - zI)^{-1}B_+(A_- - zI)^{-1}[I - B_+(A_+ - zI)^{-1}], \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (3.45)
\end{aligned}$$

Thus, relying on Lemma 3.6 and (3.36) it suffices to prove that

$$\text{s-lim}_{n \rightarrow \infty} (A_{+,n} - zI)^{-1} = (A_+ - zI)^{-1}, \quad z \in \mathbb{C} \setminus \mathbb{R}, \quad (3.46)$$

but this immediately follows from

$$(A_{+,n} - zI)^{-1} = [I + (A_- - zI)^{-1}\chi_n(A_-)B_+\chi_n(A_-)]^{-1}(A_- - zI)^{-1}, \quad (3.47)$$

$$(A_+ - zI)^{-1} = [I + (A_- - zI)^{-1}B_+]^{-1}(A_- - zI)^{-1}, \quad (3.48)$$

employing the fact that strong convergence for a sequence of bounded operators is equivalent to strong resolvent convergence, initially, for $|\text{Im}(z)|$ sufficiently large, and subsequently, for all $z \in \mathbb{C} \setminus \mathbb{R}$ by analytic continuation with respect to z . \square

Next, going beyond the approximation $A_{+,n}$ of A_+ , we now introduce the following path $\{A_+(s)\}_{s \in [0,1]}$, where

$$A_+(s) = A_- + \tilde{\chi}_s(A_-)B_+\tilde{\chi}_s(A_-), \quad \text{dom}(A_+(s)) = \text{dom}(A_-), \quad s \in [0,1], \quad (3.49)$$

$$\tilde{\chi}_s(\nu) = [(1-s)\nu^2 + 1]^{-1/2}, \quad \nu \in \mathbb{R}, \quad s \in [0,1], \quad (3.50)$$

in particular,

$$A_+(0) = A_{+,1} \text{ (cf. (3.40) with } n=1 \text{) and } A_+(1) = A_+. \quad (3.51)$$

Moreover, in complete analogy to (3.44), the family $A_+(s)$ depends continuously on $s \in [0,1]$ with respect to the metric

$$d(A, A') = \|(A - iI)^{-1} - (A' - iI)^{-1}\|_{\mathcal{B}_1(\mathcal{H})} \quad (3.52)$$

for A, A' in the set of self-adjoint operators which are resolvent comparable with respect to A_- (equivalently, A_+), that is, A, A' satisfy for some (and hence for all) $\zeta \in \mathbb{C} \setminus \mathbb{R}$,

$$[(A - \zeta I)^{-1} - (A_- - \zeta I)^{-1}], [(A' - \zeta I)^{-1} - (A_- - \zeta I)^{-1}] \in \mathcal{B}_1(\mathcal{H}). \quad (3.53)$$

Thus, the hypotheses of [28, Lemma 8.7.5] are satisfied and hence one obtains the following result:

Theorem 3.8. *Assume Hypotheses 3.1 and 3.2 (i) and introduce the path $A_+(s)$, $s \in [0,1]$, as in (3.49), with $A_+(0) = A_{+,1}$ (cf. (3.40) with $n=1$) and $A_+(1) = A_+$. Then for each $s \in [0,1]$, there exists a unique spectral shift function $\xi(\cdot; A_+(s), A_-)$ for the pair $(A_+(s), A_-)$, depending continuously on $s \in [0,1]$ in the $L^1(\mathbb{R}; (\nu^2 + 1)^{-1}d\nu)$ -norm, satisfying $\xi(\cdot; A_+(0), A_-) = \xi(\cdot; A_{+,1}, A_-)$, and (cf. (A.14)),*

$$2i \int_{\mathbb{R}} \frac{\xi(\lambda; A_+(s), A_-)d\lambda}{\lambda^2 + 1} = \text{tr}_{\mathcal{H}}(\ln(U_+(s)U_-^{-1})), \quad (3.54)$$

where

$$U_- = (A_- - iI)(A_- + iI)^{-1}, \quad U_+(s) = (A_+(s) - iI)(A_+(s) + iI)^{-1}, \quad s \in [0,1]. \quad (3.55)$$

In addition (cf. (A.34)),

$$\xi(\cdot; A_+(s), A_-) \in L^1(\mathbb{R}; d\nu), \quad s \in [0,1]. \quad (3.56)$$

Thus, observing the equality $\chi_n(\cdot) = \tilde{\chi}_{(1-n^{-2})}(\cdot)$, Theorem 3.8 implies

$$\begin{aligned} 0 &= \lim_{s \uparrow 1} \|\xi(\cdot; A_+(s), A_-) - \xi(\cdot; A_+, A_-)\|_{L^1(\mathbb{R}; (\nu^2+1)^{-1} d\nu)} \\ &= \lim_{s \uparrow 1} \|\xi(\cdot; A_+(s), A_-) - \xi(\cdot; A_+(1), A_-)\|_{L^1(\mathbb{R}; (\nu^2+1)^{-1} d\nu)} \end{aligned} \quad (3.57)$$

$$= \lim_{n \rightarrow \infty} \|\xi(\cdot; A_{+,n}, A_-) - \xi(\cdot; A_+(1), A_-)\|_{L^1(\mathbb{R}; (\nu^2+1)^{-1} d\nu)}. \quad (3.58)$$

In particular, a subsequence of $\{\xi(\cdot; A_{+,n}, A_-)\}_{n \in \mathbb{N}}$ converges pointwise a.e. to $\xi(\cdot; A_+, A_-)$ as $n \rightarrow \infty$. Thus, the sequence $\xi(\cdot; A_{+,n}, A_-) \in L^1(\mathbb{R}; d\nu)$, $n \in \mathbb{N}$, naturally enforces a choice for the open constant inherent in $\xi(\cdot; A_+, A_-)$ determined by $\xi(\cdot; A_{+,1}, A_-) = \xi(\cdot; A_+(0), A_-)$, which will henceforth be adopted for the remainder of this paper.

We continue with an elementary but useful consequence of Theorem 3.8.

Corollary 3.9. *Assume Hypotheses 3.1 and 3.2 (i) and suppose that $f \in L^\infty(\mathbb{R})$. Then*

$$\lim_{n \rightarrow \infty} \|\xi(\cdot; A_{+,n}, A_-)f - \xi(\cdot; A_+, A_-)f\|_{L^1(\mathbb{R}; (\nu^2+1)^{-1} d\nu)} = 0, \quad (3.59)$$

in particular,

$$\lim_{n \rightarrow \infty} \int_{\mathbb{R}} \xi(\nu; A_{+,n}, A_-) d\nu g(\nu) = \int_{\mathbb{R}} \xi(\nu; A_+, A_-) d\nu g(\nu) \quad (3.60)$$

for all $g \in L^\infty(\mathbb{R})$ such that $\text{ess. sup}_{\nu \in \mathbb{R}} |(\nu^2 + 1)g(\nu)| < \infty$.

Proof. Relation (3.59) is clear from Theorem 3.8 and

$$\begin{aligned} &\|\xi(\cdot; A_{+,n}, A_-)f - \xi(\cdot; A_+, A_-)f\|_{L^1(\mathbb{R}; (\nu^2+1)^{-1} d\nu)} \\ &\leq \|f\|_{L^\infty(\mathbb{R})} \|\xi(\cdot; A_{+,n}, A_-) - \xi(\cdot; A_+, A_-)\|_{L^1(\mathbb{R}; (\nu^2+1)^{-1} d\nu)}, \end{aligned} \quad (3.61)$$

and (3.60) is obvious from (3.59) and decomposing the (complex) measures

$$\xi(\nu; A_{+,n}, A_-) d\nu g(\nu) \quad \text{and} \quad \xi(\nu; A_+, A_-) d\nu g(\nu) \quad (3.62)$$

into

$$(\nu^2 + 1)^{-1} \xi(\nu; A_{+,n}, A_-) d\nu (\nu^2 + 1)g(\nu) \quad \text{and} \quad (\nu^2 + 1)^{-1} \xi(\nu; A_+, A_-) d\nu (\nu^2 + 1)g(\nu). \quad (3.63)$$

□

At this point we introduce one more assumption regarding $B'_n(t)$, $n \in \mathbb{N}$, $t \in \mathbb{R}$, and for convenience now collect all our hypotheses at one place:

Hypothesis 3.10. *Suppose \mathcal{H} is a complex, separable Hilbert space.*

(i) *Assume A_- is self-adjoint on $\text{dom}(A_-) \subseteq \mathcal{H}$.*

(ii) *Suppose we have a family of bounded operators $\{B(t)\}_{t \in \mathbb{R}} \subset \mathcal{B}(\mathcal{H})$, continuously differentiable in norm on \mathbb{R} such that*

$$\|B'(\cdot)\|_{\mathcal{B}(\mathcal{H})} \in L^1(\mathbb{R}; dt). \quad (3.64)$$

(iii) *Suppose that $|B_+|^{1/2}(A_- - z_0 I)^{-1} \in \mathcal{B}_2(\mathcal{H})$ for some (and hence for all) $z_0 \in \rho(A_-)$. (Here $B_+ = \text{n-lim}_{t \rightarrow +\infty} B(t)$.)*

(iv) *Assume that $\sup_{t \in \mathbb{R}} \|B'(t)\|_{\mathcal{B}(\mathcal{H})} < \infty$.*

(v) *Suppose that $A_- B$ is bounded with respect to H_0 with bound strictly less than one, that is, there exists $0 \leq a < 1$ and $b \in (0, \infty)$ such that*

$$\|A_- B f\|_{L^2(\mathbb{R}; \mathcal{H})} \leq a \|H_0 f\|_{L^2(\mathbb{R}; \mathcal{H})} + b \|f\|_{L^2(\mathbb{R}; \mathcal{H})}, \quad f \in \text{dom}(H_0). \quad (3.65)$$

(vi) Assume that for some (and hence for all) $z_0 \in \mathbb{C} \setminus \mathbb{R}$,

$$|\mathbf{B}'|^{1/2}(\mathbf{H}_0 - z_0 \mathbf{I})^{-1} \in \mathcal{B}_2(L^2(\mathbb{R}; \mathcal{H})) \quad (3.66)$$

(vii) Assume that

$$\mathbf{B}'_n(t) \in \mathcal{B}_1(\mathcal{H}), \quad \|B'_n(\cdot)\|_{\mathcal{B}_1(\mathcal{H})} \in L^1(\mathbb{R}; dt), \quad n \in \mathbb{N}, t \in \mathbb{R}. \quad (3.67)$$

Remark 3.11. We note that the final two assumptions (vi) and (vii) in Hypothesis 3.10 can be derived from the stronger condition

$$|B'(t)|^{1/2}(|A_-| + I)^{-1} \in \mathcal{B}_2(\mathcal{H}), \quad \| |B'(\cdot)|^{1/2}(|A_-| + I)^{-1} \|_{\mathcal{B}_2(\mathcal{H})} \in L^2(\mathbb{R}; dt). \quad (3.68)$$

Indeed, repeating the argument in [14, Lemma 4.6], one can obtain the inclusion $|\mathbf{B}'|^{1/2}(\mathbf{H}_0 - z_0 \mathbf{I})^{-1} \in \mathcal{B}_2(L^2(\mathbb{R}; \mathcal{H}))$. In addition, with regard to Hypothesis 3.10 (vii), one obtains

$$\begin{aligned} \int_{\mathbb{R}} dt \|\mathbf{B}'_n(t)\|_{\mathcal{B}_1(\mathcal{H})} &= \int_{\mathbb{R}} dt \|\chi_n(A_-)B'(t)\chi_n(A_-)\|_{\mathcal{B}_1(\mathcal{H})} \\ &\leq \int_{\mathbb{R}} dt \|\overline{\chi_n(A_-)(|A_-| + I)}\|_{\mathcal{B}(\mathcal{H})} \| |B'(t)|^{1/2}(|A_-| + I)^{-1} \|_{\mathcal{B}_2(\mathcal{H})}^2 \\ &\quad \times \|(|A_-| + I)\chi_n(A_-)\|_{\mathcal{B}(\mathcal{H})}, \end{aligned} \quad (3.69)$$

and since $\overline{\chi_n(A_-)(|A_-| + I)}, (|A_-| + I)\chi_n(A_-) \in \mathcal{B}(\mathcal{H})$, one infers that

$$\int_{\mathbb{R}} dt \|\mathbf{B}'_n(t)\|_{\mathcal{B}_1(\mathcal{H})} \leq \int_{\mathbb{R}} dt \| |B'(t)|^{1/2}(|A_-| + I)^{-1} \|_{\mathcal{B}_2(\mathcal{H})}^2 < \infty. \quad (3.70)$$

◇

Assuming Hypothesis 3.10 from now on, one obtains the decompositions,

$$\begin{aligned} \mathbf{H}_{j,n} &= \frac{d^2}{dt^2} + \mathbf{A}_n^2 + (-1)^j \mathbf{A}'_n \\ &= \mathbf{H}_0 + \mathbf{B}_n \mathbf{A}_- + \mathbf{A}_- \mathbf{B}_n + \mathbf{B}_n^2 + (-1)^j \mathbf{B}'_n, \\ \text{dom}(\mathbf{H}_{j,n}) &= \text{dom}(\mathbf{H}_0) = W^{2,2}(\mathbb{R}; \mathcal{H}), \quad n \in \mathbb{N}, j = 1, 2, \end{aligned} \quad (3.71)$$

with

$$\mathbf{B}_n = \chi_n(\mathbf{A}_-) \mathbf{B} \chi_n(\mathbf{A}_-), \quad \mathbf{B}'_n = \chi_n(\mathbf{A}_-) \mathbf{B}' \chi_n(\mathbf{A}_-), \quad n \in \mathbb{N}. \quad (3.72)$$

Lemma 3.12. Assume Hypothesis 3.10 and let $z \in \mathbb{C} \setminus [0, \infty)$. Then the following assertions hold:

(i) The operators $\mathbf{H}_{j,n}$ converge to \mathbf{H}_j , $j = 1, 2$, in the strong resolvent sense,

$$\text{s-lim}_{n \rightarrow \infty} (\mathbf{H}_{j,n} - z \mathbf{I})^{-1} = (\mathbf{H}_j - z \mathbf{I})^{-1}, \quad j = 1, 2. \quad (3.73)$$

(ii) The operators

$$\begin{aligned} &(\mathbf{H}_0 - z \mathbf{I})(\mathbf{H}_{j,n} - z \mathbf{I})^{-1}, \quad n \in \mathbb{N}, j = 1, 2, \\ &\overline{(\mathbf{H}_{j,n} - z \mathbf{I})^{-1}(\mathbf{H}_0 - z \mathbf{I})} = [(\mathbf{H}_0 - \bar{z} \mathbf{I})(\mathbf{H}_{j,n} - \bar{z} \mathbf{I})^{-1}]^*, \quad n \in \mathbb{N}, j = 1, 2, \end{aligned} \quad (3.74)$$

are uniformly bounded with respect to $n \in \mathbb{N}$, that is, there exists $C \in (0, \infty)$ such that

$$\|(\mathbf{H}_0 - z \mathbf{I})(\mathbf{H}_{j,n} - z \mathbf{I})^{-1}\|_{\mathcal{B}(L^2(\mathbb{R}; \mathcal{H}))} \leq C, \quad n \in \mathbb{N}, j = 1, 2. \quad (3.75)$$

In addition,

$$\text{s-lim}_{n \rightarrow \infty} \overline{(\mathbf{H}_{j,n} - z\mathbf{I})^{-1}(\mathbf{H}_0 - z\mathbf{I})} = \overline{(\mathbf{H}_j - z\mathbf{I})^{-1}(\mathbf{H}_0 - z\mathbf{I})}, \quad j = 1, 2, \quad (3.76)$$

$$\text{s-lim}_{n \rightarrow \infty} (\mathbf{H}_0 - z\mathbf{I})(\mathbf{H}_{j,n} - z\mathbf{I})^{-1} = (\mathbf{H}_0 - z\mathbf{I})(\mathbf{H}_j - z\mathbf{I})^{-1}, \quad j = 1, 2. \quad (3.77)$$

Proof. Since the proof for the operators $\mathbf{H}_{2,n}, \mathbf{H}_2$ is a verbatim repetition of the proof for $\mathbf{H}_{1,n}, \mathbf{H}_1$, we exclusively focus on the latter.

(i) By (3.27) and the analogous equation for $\mathbf{H}_{1,n}$, the operators \mathbf{H}_1 and $\mathbf{H}_{1,n}$ have a common core $\text{dom}(\mathbf{H}_1)$. Since

$$\mathbf{B}' - \mathbf{B}'_n = \mathbf{B}' - \chi_n(\mathbf{A}_-)\mathbf{B}'\chi_n(\mathbf{A}_-) = (\mathbf{I} - \chi_n(\mathbf{A}_-))\mathbf{B}' + \chi_n(\mathbf{A}_-)\mathbf{B}'(\mathbf{I} - \chi_n(\mathbf{A}_-)), \quad (3.78)$$

and \mathbf{B}' is a bounded operator, the convergence

$$\text{s-lim}_{n \rightarrow \infty} \mathbf{B}'_n = \mathbf{B}' \quad (3.79)$$

holds, employing

$$\text{s-lim}_{n \rightarrow \infty} \chi_n(\mathbf{A}_-) = \mathbf{I} \quad (3.80)$$

(applying the spectral theorem, see also (3.36)). Arguing analogously, one also obtains that

$$\text{s-lim}_{n \rightarrow \infty} \mathbf{B}_n = \mathbf{B}. \quad (3.81)$$

Next, rewriting

$$\begin{aligned} \mathbf{B}^2 - \mathbf{B}_n^2 &= \mathbf{B}^2 - \chi_n(\mathbf{A}_-)\mathbf{B}\chi_n(\mathbf{A}_-)\mathbf{B}\chi_n(\mathbf{A}_-) \\ &= (\mathbf{I} - \chi_n(\mathbf{A}_-))\mathbf{B}^2 + \chi_n(\mathbf{A}_-)\mathbf{B}(\mathbf{B}(\mathbf{I} - \chi_n(\mathbf{A}_-)) + (\mathbf{I} - \chi_n(\mathbf{A}_-))\mathbf{B}\chi_n(\mathbf{A}_-)), \end{aligned} \quad (3.82)$$

one also obtains

$$\text{s-lim}_{n \rightarrow \infty} \mathbf{B}_n^2 = \mathbf{B}^2. \quad (3.83)$$

Thus, it remains to show that for all $f \in \text{dom}(\mathbf{H}_1)$, $\text{s-lim}_{n \rightarrow \infty} \mathbf{B}_n \mathbf{A}_- f = \mathbf{B} \mathbf{A}_- f$ and $\text{s-lim}_{n \rightarrow \infty} \mathbf{A}_- \mathbf{B}_n f = \mathbf{A}_- \mathbf{B} f$. Indeed, one verifies

$$\begin{aligned} \mathbf{B} \mathbf{A}_- - \mathbf{B}_n \mathbf{A}_- &= \mathbf{B} \mathbf{A}_- - \chi_n(\mathbf{A}_-)\mathbf{B}\chi_n(\mathbf{A}_-)\mathbf{A}_- \\ &= (\mathbf{I} - \chi_n(\mathbf{A}_-))\mathbf{B} \mathbf{A}_- + \chi_n(\mathbf{A}_-)\mathbf{B}(\mathbf{I} - \chi_n(\mathbf{A}_-))\mathbf{A}_-, \end{aligned} \quad (3.84)$$

and

$$\begin{aligned} \mathbf{A}_- \mathbf{B} - \mathbf{A}_- \mathbf{B}_n &= \mathbf{A}_- \mathbf{B} - \mathbf{A}_- \chi_n(\mathbf{A}_-)\mathbf{B}\chi_n(\mathbf{A}_-) \\ &= (\mathbf{I} - \chi_n(\mathbf{A}_-))\mathbf{A}_- \mathbf{B} + \mathbf{A}_- \chi_n(\mathbf{A}_-)\mathbf{B}(\mathbf{I} - \chi_n(\mathbf{A}_-)), \end{aligned} \quad (3.85)$$

implying the required convergence. Consequently,

$$\text{s-lim}_{n \rightarrow \infty} \mathbf{H}_{1,n} f = \mathbf{H}_1 f, \quad f \in \text{dom}(\mathbf{H}_1). \quad (3.86)$$

Since $\mathbf{H}_{1,n}$ and \mathbf{H}_1 are self-adjoint operators with a common core, [23, Theorem VIII.25] (see also [26, Theorem 9.16]) implies that $\mathbf{H}_{1,n}$ converges to \mathbf{H}_1 in the strong resolvent sense.

(ii) Fix $z \in \mathbb{C} \setminus [0, \infty)$. First, one observes that

$$\overline{(\mathbf{H}_{1,n} - z\mathbf{I})^{-1}(\mathbf{H}_0 - z\mathbf{I})} = [(\mathbf{H}_0 - \bar{z}\mathbf{I})(\mathbf{H}_{1,n} - \bar{z}\mathbf{I})^{-1}]^*. \quad (3.87)$$

Using the standard resolvent identity one obtains

$$(\mathbf{H}_{1,n} - z\mathbf{I})^{-1} - (\mathbf{H}_0 - z\mathbf{I})^{-1} = -(\mathbf{H}_{1,n} - z\mathbf{I})^{-1}[(\mathbf{H}_1 - \mathbf{H}_0)(\mathbf{H}_0 - z\mathbf{I})^{-1}], \quad (3.88)$$

and hence concludes,

$$\begin{aligned} & (\mathbf{H}_0 - z\mathbf{I})(\mathbf{H}_{1,n} - z\mathbf{I})^{-1} \\ &= [\mathbf{I} + (\mathbf{A}_-\mathbf{B}_n + \mathbf{B}_n\mathbf{A}_- + \mathbf{B}_n^2 - \mathbf{B}_n')(\mathbf{H}_0 - z\mathbf{I})^{-1}]^{-1}. \end{aligned} \quad (3.89)$$

Because of (3.79) and (3.83), it suffices to focus on the terms $\mathbf{B}_n\mathbf{A}_-$ and $\mathbf{A}_-\mathbf{B}_n$. As in (3.25) one estimates

$$\begin{aligned} & \|\mathbf{B}_n\mathbf{A}_-(\mathbf{H}_0 - z\mathbf{I})^{-1}\|_{\mathcal{B}(L^2(\mathbb{R};\mathcal{H}))} \leq |\operatorname{Im}(z)|^{-1/2} \|\mathbf{B}\|_{\mathcal{B}(L^2(\mathbb{R};\mathcal{H}))} \\ & \times \|\mathbf{A}_-(\mathbf{H}_0 - z\mathbf{I})^{-1/2}\|_{\mathcal{B}(L^2(\mathbb{R};\mathcal{H}))}, \quad n \in \mathbb{N}, \end{aligned} \quad (3.90)$$

employing $\mathbf{B}_n = \chi_n(\mathbf{A}_-)\mathbf{B}\chi_n(\mathbf{A}_-)$, $\|\chi_n(\mathbf{A}_-)\|_{\mathcal{B}(L^2(\mathbb{R};\mathcal{H}))} \leq 1$, $n \in \mathbb{N}$, and commutativity of $\chi_n(\mathbf{A}_-)$ and $(\mathbf{H}_0 - z\mathbf{I})^{-1}$, that is,

$$[\chi_n(\mathbf{A}_-), (\mathbf{H}_0 - z\mathbf{I})^{-1}] = 0, \quad n \in \mathbb{N}. \quad (3.91)$$

Similarly, utilizing the estimate (3.26), one concludes

$$\begin{aligned} & \|\mathbf{A}_-\mathbf{B}_n(\mathbf{H}_0 - z\mathbf{I})^{-1}\|_{\mathcal{B}(L^2(\mathbb{R};\mathcal{H}))} \leq \|\mathbf{A}_-\mathbf{B}(\mathbf{H}_0 - z\mathbf{I})^{-1}\|_{\mathcal{B}(L^2(\mathbb{R};\mathcal{H}))} \\ & \leq a < 1, \quad n \in \mathbb{N}, \end{aligned} \quad (3.92)$$

for $0 < |\operatorname{Im}(z)|$ sufficiently large. Thus, choosing $0 < |\operatorname{Im}(z)|$ sufficiently large, the operator in (3.89) is uniformly bounded in norm, proving (3.75). In fact, using $\mathbf{B}_n = \chi_n(\mathbf{A}_-)\mathbf{B}\chi_n(\mathbf{A}_-)$ again, and repeatedly employing commutativity of $\chi_n(\mathbf{A}_-)$ and $(\mathbf{H}_0 - z\mathbf{I})^{-1}$, (3.89) also proves the strong convergence of $(\mathbf{A}_-\mathbf{B}_n + \mathbf{B}_n\mathbf{A}_- + \mathbf{B}_n^2 - \mathbf{B}_n')(\mathbf{H}_0 - z\mathbf{I})^{-1}$ to $(\mathbf{A}_-\mathbf{B} + \mathbf{B}\mathbf{A}_- + \mathbf{B}^2 - \mathbf{B}')$ as $n \rightarrow \infty$ for $0 < |\operatorname{Im}(z)|$ sufficiently large. Using the fact that strong convergence of a sequence of uniformly bounded operators is equivalent to strong resolvent convergence of the sequence, one obtains the asserted strong convergence in (3.77) for $0 < |\operatorname{Im}(z)|$ sufficiently large. An application of (3.73) together with the bound (3.75) permits one to extend this to all $z \in \mathbb{C} \setminus [0, \infty)$, completing the proof of (3.77).

Finally, to prove (3.76) it suffices to combine the strong resolvent convergence in (3.73), the uniform boundedness in (3.75) with equality (3.74), the strong convergence

$$\operatorname{s-lim}_{n \rightarrow \infty} (\mathbf{H}_{1,n} - z\mathbf{I})^{-1}(\mathbf{H}_0 - z\mathbf{I})f = (\mathbf{H}_1 - z\mathbf{I})^{-1}(\mathbf{H}_0 - z\mathbf{I})f, \quad f \in \operatorname{dom}(\mathbf{H}_0), \quad (3.93)$$

and the fact that $\operatorname{dom}(\mathbf{H}_0)$ is dense in $L^2(\mathbb{R};\mathcal{H})$. Here we used that uniformly bounded sequences of bounded operators in a Hilbert space converge strongly if they converge pointwise on a dense subset of the Hilbert space. \square

Next, we recall that

$$\mathbf{H}_2 - \mathbf{H}_1 = 2\mathbf{B}', \quad (3.94)$$

$$\mathbf{H}_{2,n} - \mathbf{H}_{1,n} = 2\mathbf{B}'_n = 2\chi_n(\mathbf{A}_-)\mathbf{B}'\chi_n(\mathbf{A}_-), \quad n \in \mathbb{N}, \quad (3.95)$$

and in analogy to (3.28)–(3.31) one concludes that

$$[(\mathbf{H}_{2,n} - z\mathbf{I})^{-1} - (\mathbf{H}_{1,n} - z\mathbf{I})^{-1}] \in \mathcal{B}_1(L^2(\mathbb{R};\mathcal{H})), \quad n \in \mathbb{N}, \quad z \in \mathbb{C} \setminus [0, \infty), \quad (3.96)$$

since

$$\begin{aligned} & (\mathbf{H}_{2,n} - z\mathbf{I})^{-1} - (\mathbf{H}_{1,n} - z\mathbf{I})^{-1} = -[(\mathbf{H}_{1,n} - z\mathbf{I})^{-1}(\mathbf{H}_0 - z\mathbf{I})]\chi_n(\mathbf{A}_-) \\ & \times [(\mathbf{H}_0 - z\mathbf{I})^{-1}2\mathbf{B}'(\mathbf{H}_0 - z\mathbf{I})^{-1}]\chi_n(\mathbf{A}_-)[(\mathbf{H}_0 - z\mathbf{I})(\mathbf{H}_{2,n} - z\mathbf{I})^{-1}], \end{aligned} \quad (3.97)$$

$$n \in \mathbb{N}, z \in \mathbb{C} \setminus [0, \infty),$$

again employing commutativity of $\chi_n(\mathbf{A}_-)$ and $(\mathbf{H}_0 - z\mathbf{I})^{-1}$ (cf. (3.91)).

Finally, we proceed to some crucial convergence results to be used in Section 4.

Theorem 3.13. *Assume Hypothesis 3.10. Then*

$$\begin{aligned} \lim_{n \rightarrow \infty} \left\| [(\mathbf{H}_{2,n} - z\mathbf{I})^{-1} - (\mathbf{H}_{1,n} - z\mathbf{I})^{-1}] \right. \\ \left. - [(\mathbf{H}_2 - z\mathbf{I})^{-1} - (\mathbf{H}_1 - z\mathbf{I})^{-1}] \right\|_{\mathcal{B}_1(L^2(\mathbb{R}; \mathcal{H}))} = 0, \quad z \in \mathbb{C} \setminus \mathbb{R}. \end{aligned} \quad (3.98)$$

Proof. Equations (3.28) and (3.95) yield

$$\begin{aligned} & [(\mathbf{H}_{2,n} - z\mathbf{I})^{-1} - (\mathbf{H}_{1,n} - z\mathbf{I})^{-1}] - [(\mathbf{H}_2 - z\mathbf{I})^{-1} - (\mathbf{H}_1 - z\mathbf{I})^{-1}] \\ &= -2(\mathbf{H}_{2,n} - z\mathbf{I})^{-1} \mathbf{B}'_n (\mathbf{H}_{1,n} - z\mathbf{I})^{-1} + 2(\mathbf{H}_2 - z\mathbf{I})^{-1} \mathbf{B}' (\mathbf{H}_1 - z\mathbf{I})^{-1} \\ &= -2 \overline{[(\mathbf{H}_{2,n} - z\mathbf{I})^{-1} (\mathbf{H}_0 - z\mathbf{I})]} \\ &\quad \times \{ \chi_n(\mathbf{A}_-) (\mathbf{H}_0 - z\mathbf{I})^{-1} \mathbf{B}' (\mathbf{H}_0 - z\mathbf{I})^{-1} \chi_n(\mathbf{A}_-) \} \\ &\quad \times [(\mathbf{H}_0 - z\mathbf{I}) (\mathbf{H}_{1,n} - z\mathbf{I})^{-1}], \end{aligned} \quad (3.99)$$

$$\begin{aligned} &+ 2 \overline{[(\mathbf{H}_2 - z\mathbf{I})^{-1} (\mathbf{H}_0 - z\mathbf{I})]} \\ &\quad \times \{ (\mathbf{H}_0 - z\mathbf{I})^{-1} \mathbf{B}' (\mathbf{H}_0 - z\mathbf{I})^{-1} \} \\ &\quad \times [(\mathbf{H}_0 - z\mathbf{I}) (\mathbf{H}_1 - z\mathbf{I})^{-1}], \quad z \in \mathbb{C} \setminus [0, \infty). \end{aligned} \quad (3.100)$$

By Lemma 3.6 and (3.66), the term $\{ \chi_n(\mathbf{A}_-) (\mathbf{H}_0 - z\mathbf{I})^{-1} \mathbf{B}' (\mathbf{H}_0 - z\mathbf{I})^{-1} \chi_n(\mathbf{A}_-) \}$ converges to $\{ (\mathbf{H}_0 - z\mathbf{I})^{-1} \mathbf{B}' (\mathbf{H}_0 - z\mathbf{I})^{-1} \}$ in $\mathcal{B}_1(L^2(\mathbb{R}; \mathcal{H}))$ -norm as $n \rightarrow \infty$. Another application of Lemma 3.6 proves (3.98) since by Lemma 3.12 (ii), for $z \in \mathbb{C} \setminus [0, \infty)$, one has

$$\text{s-lim}_{n \rightarrow \infty} [(\mathbf{H}_0 - z\mathbf{I}) (\mathbf{H}_{1,n} - z\mathbf{I})^{-1}] = [(\mathbf{H}_0 - z\mathbf{I}) (\mathbf{H}_1 - z\mathbf{I})^{-1}], \quad (3.101)$$

$$\text{s-lim}_{n \rightarrow \infty} \overline{[(\mathbf{H}_{2,n} - z\mathbf{I})^{-1} (\mathbf{H}_0 - z\mathbf{I})]} = \overline{[(\mathbf{H}_2 - z\mathbf{I})^{-1} (\mathbf{H}_0 - z\mathbf{I})]}. \quad (3.102)$$

□

Theorem 3.14. *Assume Hypothesis 3.10 and let $z, z' \in \mathbb{C} \setminus [0, \infty)$. Then*

$$\lim_{n \rightarrow \infty} \left\| \mathbf{B}'_n (\mathbf{H}_{j,n} - z\mathbf{I})^{-1} - \mathbf{B}' (\mathbf{H}_j - z\mathbf{I})^{-1} \right\|_{\mathcal{B}_2(L^2(\mathbb{R}; \mathcal{H}))} = 0, \quad j = 1, 2, \quad (3.103)$$

$$\begin{aligned} \lim_{n \rightarrow \infty} \left\| (\mathbf{H}_{j,n} - z\mathbf{I})^{-1} 2\mathbf{B}'_n (\mathbf{H}_{j,n} - z'\mathbf{I})^{-1} \right. \\ \left. - (\mathbf{H}_j - z\mathbf{I})^{-1} 2\mathbf{B}' (\mathbf{H}_j - z'\mathbf{I})^{-1} \right\|_{\mathcal{B}_1(L^2(\mathbb{R}; \mathcal{H}))} = 0, \quad j = 1, 2. \end{aligned} \quad (3.104)$$

Proof. Fix $z, z' \in \mathbb{C} \setminus [0, \infty)$. To prove (3.103) one writes

$$\begin{aligned} \mathbf{B}'_n (\mathbf{H}_{j,n} - z\mathbf{I})^{-1} &= \chi_n(\mathbf{A}_-) [\mathbf{B}' (\mathbf{H}_0 - z\mathbf{I})^{-1}] \\ &\quad \times \chi_n(\mathbf{A}_-) [(\mathbf{H}_0 - z\mathbf{I}) (\mathbf{H}_{j,n} - z\mathbf{I})^{-1}], \quad j = 1, 2, n \in \mathbb{N}, \end{aligned} \quad (3.105)$$

employing once again commutativity of $\chi_n(\mathbf{A}_-)$ and $(\mathbf{H}_0 - z\mathbf{I})^{-1}$ (cf. (3.91)). Since

$$\mathbf{B}' (\mathbf{H}_0 - z\mathbf{I})^{-1} = [U_{\mathbf{B}'} |\mathbf{B}'|^{1/2}] |\mathbf{B}'|^{1/2} (\mathbf{H}_0 - z\mathbf{I})^{-1} \in \mathcal{B}_2(L^2(\mathbb{R}; \mathcal{H})), \quad (3.106)$$

by (3.24) and the polar decomposition $\mathbf{B}' = U_{\mathbf{B}'} |\mathbf{B}'|$ of $\mathbf{B}' \in \mathcal{B}(L^2(\mathbb{R}; \mathcal{H}))$ (cf. (3.21)). Thus, (3.103) is a consequence of Lemma 3.6 combined with (3.80) and (3.77).

Relation (3.104) follows along exactly the same lines upon decomposing

$$\begin{aligned} (\mathbf{H}_{j,n} - z\mathbf{I})^{-1} 2\mathbf{B}'_n(\mathbf{H}_{j,n} - z'\mathbf{I})^{-1} &= [(\overline{\mathbf{H}_{j,n} - z\mathbf{I}})^{-1}(\mathbf{H}_0 - z\mathbf{I})] \chi_n(\mathbf{A}_-) \\ &\times [(\mathbf{H}_0 - z\mathbf{I})^{-1} 2\mathbf{B}'(\mathbf{H}_0 - z'\mathbf{I})^{-1}] \chi_n(\mathbf{A}_-) [(\mathbf{H}_0 - z'\mathbf{I})(\mathbf{H}_{j,n} - z'\mathbf{I})^{-1}], \\ &j = 1, 2, n \in \mathbb{N}, \end{aligned} \quad (3.107)$$

applying once more Lemma 3.6, (3.76), (3.77), (3.80), (3.91), and

$$\begin{aligned} (\mathbf{H}_0 - z\mathbf{I})^{-1} \mathbf{B}'(\mathbf{H}_0 - z'\mathbf{I})^{-1} &= [(\mathbf{H}_0 - z\mathbf{I})^{-1} |\mathbf{B}'|^{1/2}] \operatorname{sgn}(\mathbf{B}') \\ &\times [|\mathbf{B}'|^{1/2}(\mathbf{H}_0 - z'\mathbf{I})^{-1}] \in \mathcal{B}_1(L^2(\mathbb{R}; \mathcal{H})), \end{aligned} \quad (3.108)$$

since $|\mathbf{B}'|^{1/2}(\mathbf{H}_0 - z\mathbf{I})^{-1} \in \mathcal{B}_2(L^2(\mathbb{R}; \mathcal{H}))$ by (3.24). \square

4. COMPUTING $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ IN TERMS OF $\xi(\cdot; A_+, A_-)$

Given the results of Section 3 and Appendix A, we now determine $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ in terms of $\xi(\cdot; A_+, A_-)$. This represents one of the principal results of this paper.

Theorem 4.1. *Assume Hypothesis 3.10. Then,*

$$\begin{aligned} &\int_{[0, \infty)} \xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}] \\ &= \int_{\mathbb{R}} \xi(\nu; A_+, A_-) d\nu [(\nu^2 - z)^{-1/2} - (\nu^2 - z_0)^{-1/2}], \quad z, z_0 \in \mathbb{C} \setminus [0, \infty). \end{aligned} \quad (4.1)$$

Moreover,

$$\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) = \frac{1}{\pi} \int_{-\lambda^{1/2}}^{\lambda^{1/2}} \frac{\xi(\nu; A_+, A_-) d\nu}{(\lambda - \nu^2)^{1/2}} \text{ for a.e. } \lambda > 0. \quad (4.2)$$

Proof. Due to relation (3.67), [8] and [22] apply and one concludes the approximate trace formula,

$$\begin{aligned} &\operatorname{tr}_{L^2(\mathbb{R}; \mathcal{H})} ((\mathbf{H}_{2,n} - z\mathbf{I})^{-1} - (\mathbf{H}_{1,n} - z\mathbf{I})^{-1}) \\ &= \frac{1}{2z} \operatorname{tr}_{\mathcal{H}} (g_z(A_{+,n}) - g_z(A_-)), \quad n \in \mathbb{N}, z \in \mathbb{C} \setminus [0, \infty), \end{aligned} \quad (4.3)$$

with

$$g_z(x) = x(x^2 - z)^{-1/2}, \quad z \in \mathbb{C} \setminus [0, \infty), x \in \mathbb{R}. \quad (4.4)$$

Relation (4.3) and the Krein–Lifshitz trace formula yield

$$\int_{[0, \infty)} \frac{\xi(\lambda; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) d\lambda}{(\lambda - z)^2} = \frac{1}{2} \int_{\mathbb{R}} \frac{\xi(\nu; A_{+,n}, A_-) d\nu}{(\nu^2 - z)^{3/2}}, \quad n \in \mathbb{N}, z \in \mathbb{C} \setminus [0, \infty). \quad (4.5)$$

As shown in the course of the proof of Theorem 8.2 in [14], (4.5) implies the relation

$$\begin{aligned} &\int_{[0, \infty)} \xi(\lambda; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}] \\ &= \int_{\mathbb{R}} \xi(\nu; A_{+,n}, A_-) d\nu [(\nu^2 - z)^{-1/2} - (\nu^2 - z_0)^{-1/2}], \\ &n \in \mathbb{N}, z, z_0 \in \mathbb{C} \setminus [0, \infty). \end{aligned} \quad (4.6)$$

Combining Theorem 3.13 with the Krein–Lifshitz trace formula (3.34) (for the pair $(\mathbf{H}_2, \mathbf{H}_1)$ as well as the pairs $(\mathbf{H}_{2,n}, \mathbf{H}_{1,n})$, $n \in \mathbb{N}$), yields

$$\begin{aligned} & \lim_{n \rightarrow \infty} \int_{[0, \infty)} \frac{\xi(\lambda; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) d\lambda}{(\lambda - z)^2} \\ &= - \lim_{n \rightarrow \infty} \operatorname{tr}_{L^2(\mathbb{R}; \mathcal{H})} ((\mathbf{H}_{2,n} - z\mathbf{I})^{-1} - (\mathbf{H}_{1,n} - z\mathbf{I})^{-1}) \\ &= - \operatorname{tr}_{L^2(\mathbb{R}; \mathcal{H})} ((\mathbf{H}_2 - z\mathbf{I})^{-1} - (\mathbf{H}_1 - z\mathbf{I})^{-1}) \\ &= \int_{[0, \infty)} \frac{\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) d\lambda}{(\lambda - z)^2}, \quad z \in \mathbb{C} \setminus \mathbb{R}. \end{aligned} \quad (4.7)$$

Lemma 3.12 (i) and Theorem 3.13 imply that the pairs of self-adjoint operators $(\mathbf{H}_{2,n}, \mathbf{H}_{1,n})$, $n \in \mathbb{N}$, and $(\mathbf{H}_2, \mathbf{H}_1)$ satisfy the hypotheses (A.45), (A.51), (A.53), (A.57) (identifying the pairs $(A_n, A_{0,n})$ and (A, A_0) with the pairs $(\mathbf{H}_{2,n}, \mathbf{H}_{1,n})$ and $(\mathbf{H}_2, \mathbf{H}_1)$, respectively). Thus, an application of (A.58) to the pairs $(\mathbf{H}_{2,n}, \mathbf{H}_{1,n})$ and $(\mathbf{H}_2, \mathbf{H}_1)$ implies

$$\begin{aligned} & \lim_{n \rightarrow \infty} \int_{[0, \infty)} \xi(\lambda; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}] \\ &= \int_{[0, \infty)} \xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}], \quad z, z_0 \in \mathbb{C} \setminus \mathbb{R}. \end{aligned} \quad (4.8)$$

On the other hand, since

$$[(\nu^2 - z)^{-1/2} - (\nu^2 - z_0)^{-1/2}]_{|\nu| \rightarrow \infty} = O(|\nu|^{-3}), \quad z, z_0 \in \mathbb{C} \setminus \mathbb{R}, \quad (4.9)$$

and hence $(\nu^2 + 1)[(\nu^2 - z)^{-1/2} - (\nu^2 - z_0)^{-1/2}]$ is uniformly bounded for $\nu \in \mathbb{R}$, (3.60) yields

$$\begin{aligned} & \lim_{n \rightarrow \infty} \int_{\mathbb{R}} \xi(\nu; A_{+,n}, A_-) d\nu [(\nu^2 - z)^{-1/2} - (\nu^2 - z_0)^{-1/2}] \\ &= \int_{\mathbb{R}} \xi(\nu; A_+, A_-) d\nu [(\nu^2 - z)^{-1/2} - (\nu^2 - z_0)^{-1/2}], \quad z, z_0 \in \mathbb{C} \setminus [0, \infty). \end{aligned} \quad (4.10)$$

Thus, combining (4.6), (4.8), and (4.10) one finally obtains

$$\begin{aligned} & \int_{[0, \infty)} \xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}] \\ &= \lim_{n \rightarrow \infty} \int_{[0, \infty)} \xi(\lambda; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}] \\ &= \lim_{n \rightarrow \infty} \int_{\mathbb{R}} \xi(\nu; A_{+,n}, A_-) d\nu [(\nu^2 - z)^{-1/2} - (\nu^2 - z_0)^{-1/2}] \\ &= \int_{\mathbb{R}} \xi(\nu; A_+, A_-) d\nu [(\nu^2 - z)^{-1/2} - (\nu^2 - z_0)^{-1/2}], \quad z, z_0 \in \mathbb{C} \setminus [0, \infty), \end{aligned} \quad (4.11)$$

and hence (4.1).

Applying the Stieltjes inversion formula (see the discussion in [2] and in [26, Appendix B]) to (4.11) then proves (4.2) precisely along the lines detailed in the proof of [14, Theorem 8.2]. \square

Equation (4.2) now represents a far reaching extension of Pushnitski's formula originally obtained in [22] and considerably generalized in [14]. In particular, the relative trace class assumption employed in [14] has now been removed.

5. THE WITTEN INDEX

In this section we briefly discuss the notion of the Witten index for \mathbf{D}_A following the detailed treatment in [9]. The results of the present paper now enable us to remove the “relatively trace class perturbation assumption” in [9] as well as the Fredholm hypothesis in [14].

Definition 5.1. *Let T be a closed, linear, densely defined operator in \mathcal{H} and suppose that for some (and hence for all) $z \in \mathbb{C} \setminus [0, \infty)$,*

$$[(T^*T - zI_{\mathcal{H}})^{-1} - (TT^* - zI_{\mathcal{H}})^{-1}] \in \mathcal{B}_1(\mathcal{H}). \quad (5.1)$$

Then introducing the resolvent regularization

$$\Delta_r(T, \lambda) = (-\lambda) \operatorname{tr}_{\mathcal{H}} ((T^*T - \lambda I_{\mathcal{H}})^{-1} - (TT^* - \lambda I_{\mathcal{H}})^{-1}), \quad \lambda < 0, \quad (5.2)$$

the resolvent regularized Witten index $W_r(T)$ of T is defined by

$$W_r(T) = \lim_{\lambda \uparrow 0} \Delta_r(T, \lambda), \quad (5.3)$$

whenever this limit exists.

Here, in obvious notation, the subscript “ r ” indicates the use of the resolvent regularization (for a semigroup or heat kernel regularization we refer to [9]; the heat kernel regularization yields results consistent with the resolvent regularization, see, [9]). Before proceeding to compute the Witten index we recall the known consistency between the Fredholm and Witten index whenever T is Fredholm:

Theorem 5.2. ([4], [18].) *Suppose that T is a Fredholm operator in \mathcal{H} . If (5.1) holds, then the resolvent regularized Witten index $W_r(T)$ exists, equals the Fredholm index, $\operatorname{index}(T)$, of T , and*

$$W_r(T) = \operatorname{index}(T) = \xi(0_+; TT^*, T^*T). \quad (5.4)$$

The following result is proved in [9, Theorem 2.6] under the assumption of a relatively trace class perturbation, however, the argument can be adapted to the present setting.

Theorem 5.3. *Assume Hypothesis 3.1. Then the operator \mathbf{D}_A is Fredholm if and only if $0 \in \rho(A_+) \cap \rho(A_-)$.*

Since generally, \mathbf{D}_A is not a Fredholm operator in $L^2(\mathbb{R}; \mathcal{H})$, we now determine the resolvent regularized Witten index of \mathbf{D}_A as follows:

Theorem 5.4. *Assume Hypothesis 3.10 and assume that 0 is a right and a left Lebesgue point of $\xi(\cdot; A_+, A_-)$ (denoted by $\xi_L(0_+; A_+, A_-)$ and $\xi_L(0_-; A_+, A_-)$). Then 0 is a right Lebesgue point of $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ (denoted by $\xi_L(0_+; \mathbf{H}_2, \mathbf{H}_1)$) and $W_r(\mathbf{D}_A)$ exists and equals*

$$W_r(\mathbf{D}_A) = \xi_L(0_+; \mathbf{H}_2, \mathbf{H}_1) = [\xi_L(0_+; A_+, A_-) + \xi_L(0_-; A_+, A_-)]/2. \quad (5.5)$$

In particular, if $0 \in \rho(A_+) \cap \rho(A_-)$, then \mathbf{D}_A is Fredholm and

$$\operatorname{index}(\mathbf{D}_A) = W_r(\mathbf{D}_A) = \xi(0; A_+, A_-). \quad (5.6)$$

Proof. One can closely follow the argument as used in [9]. First, one rewrites (4.2) in the form,

$$\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) = \frac{1}{\pi} \int_0^{\lambda^{1/2}} \frac{d\nu [\xi(\nu; A_+, A_-) + \xi(-\nu; A_+, A_-)]}{(\lambda - \nu^2)^{1/2}}, \quad \lambda > 0. \quad (5.7)$$

Applying [9, Lemma 4.1 (i)] to the function $f(\nu) = [\xi(\nu, A_+, A_-) + \xi(-\nu, A_+, A_-)]$, $\nu > 0$, yields the Lebesgue point statement for $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ and also proves that $\xi_L(0_+; \mathbf{H}_2, \mathbf{H}_1) = [\xi_L(0_+; A_+, A_-) + \xi_L(0_-; A_+, A_-)]/2$.

Next, we note that combining Corollary 3.9, (4.5), (4.7) yields

$$\begin{aligned} -z \operatorname{tr}_{L^2(\mathbb{R}; \mathcal{H})} ((\mathbf{H}_2 - z\mathbf{I})^{-1} - (\mathbf{H}_1 - z\mathbf{I})^{-1}) &= z \int_{[0, \infty)} \frac{\xi(\lambda'; \mathbf{H}_2, \mathbf{H}_1) d\lambda'}{(\lambda' - z)^2} \\ &= z \lim_{n \rightarrow \infty} \int_{[0, \infty)} \frac{\xi(\lambda'; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) d\lambda'}{(\lambda' - z)^2} = \frac{z}{2} \lim_{n \rightarrow \infty} \int_{\mathbb{R}} \frac{\xi(\nu; A_{+,n}, A_-) d\nu}{(\nu^2 - z)^{3/2}}, \\ &= \frac{z}{2} \int_{\mathbb{R}} \frac{\xi(\nu; A_+, A_-) d\nu}{(\nu^2 - z)^{3/2}} \quad z \in \mathbb{C} \setminus [0, \infty). \end{aligned} \quad (5.8)$$

Thus, one obtains existence of $W_r(\mathbf{D}_A)$ (cf. (5.3)) and its equality with the expression $[\xi_L(0_+; A_+, A_-) + \xi_L(0_-; A_+, A_-)]/2$ upon applying [9, Lemma 4.2] to the last term in (5.8), with $z = \lambda < 0$.

In the case where A_{\pm} are boundedly invertible, the equality $\operatorname{index}(\mathbf{D}_A) = W_r(\mathbf{D}_A)$ immediately follows from Theorems 5.2 and 5.3. In addition, since $0 \in \rho(A_+) \cap \rho(A_-)$, the function $\xi(\cdot; A_+, A_-)$ is constant on some interval $(-\varepsilon, \varepsilon)$ for some $\varepsilon > 0$ (see, e.g., [28, p. 300]), and hence $\xi_L(0_+; A_+, A_-) = \xi_L(0_-; A_+, A_-) = \xi(0; A_+, A_-)$ yields (5.6). \square

In general (i.e., if T is not Fredholm), $W_r(T)$ is not necessarily integer-valued; in fact, it can take on any prescribed real number (cf., e.g., the analysis in [1], [4]). In this context we recall the crucial fact that $W_r(T)$ has stability properties with respect to additive perturbations analogous to the Fredholm index, as long as one replaces the familiar relative compactness assumption on the perturbation in connection with the Fredholm index, by appropriate relative trace class conditions in connection with the Witten index, as shown in [4] and [18]. In this context we also refer to [12].

6. A (1 + 1)-DIMENSIONAL EXAMPLE

In our final section we briefly illustrate our formalism in terms of a concrete (1 + 1)-dimensional example treated in great detail in [6] and [7].

Hypothesis 6.1. *Suppose the real-valued functions ϕ, θ satisfy*

$$\phi \in AC_{\text{loc}}(\mathbb{R}) \cap L^\infty(\mathbb{R}) \cap L^1(\mathbb{R}), \quad \phi' \in L^\infty(\mathbb{R}), \quad (6.1)$$

$$\theta \in AC_{\text{loc}}(\mathbb{R}) \cap L^\infty(\mathbb{R}), \quad \theta' \in L^\infty(\mathbb{R}) \cap L^1(\mathbb{R}), \quad (6.2)$$

$$\lim_{t \rightarrow \infty} \theta(t) = 1, \quad \lim_{t \rightarrow -\infty} \theta(t) = 0.$$

Given Hypothesis 6.1, we introduce the family of self-adjoint operators $A(t)$, $t \in \mathbb{R}$, in $L^2(\mathbb{R})$,

$$A(t) = -i \frac{d}{dx} + \theta(t)\phi, \quad \operatorname{dom}(A(t)) = W^{1,2}(\mathbb{R}), \quad t \in \mathbb{R}. \quad (6.3)$$

Its asymptotes as $t \rightarrow \pm\infty$ are given by

$$\mathbf{n}\text{-}\lim_{t \rightarrow \pm\infty} (A(t) - z\mathbf{I})^{-1} = (A_{\pm} - z\mathbf{I})^{-1}, \quad z \in \mathbb{C} \setminus \mathbb{R}, \quad (6.4)$$

$$A_+ = -i \frac{d}{dx} + \phi, \quad A_- = -i \frac{d}{dx}, \quad \operatorname{dom}(A_{\pm}) = W^{1,2}(\mathbb{R}). \quad (6.5)$$

(For simplicity, we adopt the abbreviation $I = I_{L^2(\mathbb{R})}$ throughout this section.)

In addition, it is convenient to introduce the family of bounded operators $B(t)$, $t \in \mathbb{R}$, in $L^2(\mathbb{R})$, where

$$B(t) = \theta(t)\phi, \quad \text{dom}(B(t)) = L^2(\mathbb{R}), \quad t \in \mathbb{R}, \quad (6.6)$$

implying

$$A(t) = A_- + B(t), \quad t \in \mathbb{R}. \quad (6.7)$$

The asymptotes of $B(t)$, $t \in \mathbb{R}$, as $t \rightarrow \pm\infty$ are then given by

$$B_+ = \text{n-lim}_{t \rightarrow +\infty} B(t) = \phi, \quad B_- = \text{n-lim}_{t \rightarrow -\infty} B(t) = 0. \quad (6.8)$$

Introducing the operator d/dt in $L^2(\mathbb{R}; dt; L^2(\mathbb{R}; dx))$ by

$$\begin{aligned} \left(\frac{d}{dt}f\right)(t) &= f'(t) \text{ for a.e. } t \in \mathbb{R}, \\ f \in \text{dom}(d/dt) &= \{g \in L^2(\mathbb{R}; dt; L^2(\mathbb{R})) \mid g \in AC_{\text{loc}}(\mathbb{R}; L^2(\mathbb{R})), \\ &\quad g' \in L^2(\mathbb{R}; dt; L^2(\mathbb{R}))\} \end{aligned} \quad (6.9)$$

$$= W^{1,2}(\mathbb{R}; dt; L^2(\mathbb{R}; dx)). \quad (6.10)$$

one now defines \mathbf{A} , \mathbf{B} , \mathbf{A}_- , $\mathbf{A}' = \mathbf{B}'$, $\mathbf{D}_{\mathbf{A}}$, \mathbf{H}_0 , and the pair $(\mathbf{H}_2, \mathbf{H}_1)$ in the Hilbert space $L^2(\mathbb{R}; dt; L^2(\mathbb{R}; dx))$ as in Section 3 and for notational simplicity we agree to identify $L^2(\mathbb{R}; dt; L^2(\mathbb{R}; dx))$ with $L^2(\mathbb{R}^2; dt dx)$ in the following. In particular, $\mathbf{D}_{\mathbf{A}}$ in $L^2(\mathbb{R}^2)$ is of the form

$$\mathbf{D}_{\mathbf{A}} = \frac{d}{dt} + \mathbf{A}, \quad \text{dom}(\mathbf{D}_{\mathbf{A}}) = W^{1,2}(\mathbb{R}^2), \quad (6.11)$$

with \mathbf{A} defined as in (1.2) identifying $\mathcal{H} = L^2(\mathbb{R})$ and $A(t)$, $t \in \mathbb{R}$, is given by (6.3).

Similarly, mimicking the approximation setup described (3.35)–(3.42) one introduces $A_n(t)$, its asymptotes $A_{\pm,n}$ as $t \rightarrow \pm\infty$, $B_n(t)$, $t \in \mathbb{R}$, $n \in \mathbb{N}$, and then also defines \mathbf{A}_n , \mathbf{B}_n , $\mathbf{A}'_n = \mathbf{B}'_n$, and the pair $(\mathbf{H}_{2,n}, \mathbf{H}_{1,n})$ in $L^2(\mathbb{R}^2; dt dx)$ as in Section 3.

As shown in [6], the assumptions on ϕ and θ made in Hypothesis 6.1 guarantee that all conditions in Hypothesis 3.10 are met and the following results can be derived:

$$[(A_+ - zI)^{-1} - (A_- - zI)^{-1}] \in \mathcal{B}_1(L^2(\mathbb{R})), \quad z \in \mathbb{C} \setminus \mathbb{R}, \quad (6.12)$$

and thus, the spectral shift function $\xi(\cdot; A_+, A_-)$ for the pair (A_+, A_-) exists and is well-defined up to an arbitrary additive real constant, satisfying

$$\xi(\cdot; A_+, A_-) \in L^1(\mathbb{R}; (\nu^2 + 1)^{-1} d\nu). \quad (6.13)$$

Introducing $\chi_n(A_-) = n(A_-^2 + n^2 I)^{-1/2}$, $n \in \mathbb{N}$, according to (3.35), the fact

$$A_{+,n} - A_- = \chi_n(A_-)B_n\chi_n(A_-) \in \mathcal{B}_1(L^2(\mathbb{R})), \quad n \in \mathbb{N}, \quad (6.14)$$

implies that the spectral shift functions $\xi(\cdot; A_{+,n}, A_-)$, $n \in \mathbb{N}$, exist and are uniquely determined by

$$\xi(\cdot; A_{+,n}, A_-) \in L^1(\mathbb{R}; d\nu), \quad n \in \mathbb{N}. \quad (6.15)$$

In fact, one can derive the expressions

$$\xi(\nu; A_{+,n}, A_-) = \pi^{-1} \text{Im}(\ln(\det_{2,L^2(\mathbb{R})}(I + \text{sgn}(\phi)|\phi|^{1/2}\chi_n(A_-)(A_- - (\nu + i0)I)^{-1}$$

$$\times \chi_n(A_-)|\phi|^{1/2})) + \frac{1}{2\pi} \frac{n^2}{\nu^2 + n^2} \int_{\mathbb{R}} dx \phi(x) \text{ for a.e. } \nu \in \mathbb{R}, \quad n \in \mathbb{N}, \quad (6.16)$$

and when studying the limit $n \rightarrow \infty$ of $\xi(\cdot; A_{+,n}, A_-)$ one can prove that

$$\lim_{n \rightarrow \infty} \xi(\nu; A_{+,n}, A_-) = \frac{1}{2\pi} \int_{\mathbb{R}} dx \phi(x), \quad \nu \in \mathbb{R}. \quad (6.17)$$

In addition, one can show that Theorem 3.8 applies and hence $\xi(\cdot; A_+, A_-)$ associated with the pair (A_+, A_-) , normalized according to (3.54) (see also the discussion in Appendix A, particularly, (A.14)) is determined via

$$\lim_{n \rightarrow \infty} \xi(\nu; A_{+,n}, A_-) = \frac{1}{2\pi} \int_{\mathbb{R}} dx \phi(x) = \xi(\nu; A_+, A_-), \quad \nu \in \mathbb{R}. \quad (6.18)$$

Thus, one obtains the remarkable fact that $\xi(\cdot; A_+, A_-)$ turns out to be constant in this example. (This phenomenon is explored and explained in detail in [6] in terms of scattering theoretic notions).

Similarly, the facts,

$$[(\mathbf{H}_2 - z\mathbf{I})^{-1} - (\mathbf{H}_1 - z\mathbf{I})^{-1}] \in \mathcal{B}_1(L^2(\mathbb{R}^2)), \quad z \in \mathbb{C} \setminus [0, \infty), \quad (6.19)$$

$$[(\mathbf{H}_{2,n} - z\mathbf{I})^{-1} - (\mathbf{H}_{1,n} - z\mathbf{I})^{-1}] \in \mathcal{B}_1(L^2(\mathbb{R}^2)), \quad n \in \mathbb{N}, \quad z \in \mathbb{C} \setminus [0, \infty), \quad (6.20)$$

show that the spectral shift functions $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ and $\xi(\cdot; \mathbf{H}_{2,n}, \mathbf{H}_{1,n})$ for the pairs $(\mathbf{H}_2, \mathbf{H}_1)$ and $(\mathbf{H}_{2,n}, \mathbf{H}_{1,n})$, $n \in \mathbb{N}$, respectively, are well-defined. In particular, they satisfy

$$\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1), \xi(\cdot; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) \in L^1(\mathbb{R}; (\lambda^2 + 1)^{-1} d\lambda), \quad n \in \mathbb{N}, \quad (6.21)$$

and since $\mathbf{H}_j \geq 0$, $\mathbf{H}_{j,n} \geq 0$, $n \in \mathbb{N}$, $j = 1, 2$, one uniquely introduces $\xi(\cdot; \mathbf{H}_2, \mathbf{H}_1)$ and $\xi(\cdot; \mathbf{H}_{2,n}, \mathbf{H}_{1,n})$, $n \in \mathbb{N}$, by requiring that

$$\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) = 0, \quad \xi(\cdot; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) = 0, \quad \lambda < 0, \quad n \in \mathbb{N}. \quad (6.22)$$

As shown in [6], one can now prove the following intimate connection between $\xi(\cdot; A_{+,n}, A_-)$ and $\xi(\cdot; \mathbf{H}_{2,n}, \mathbf{H}_{1,n})$, $n \in \mathbb{N}$, the Pushnitski-type formula,

$$\xi(\lambda; \mathbf{H}_{2,n}, \mathbf{H}_{1,n}) = \frac{1}{\pi} \int_{-\lambda^{1/2}}^{\lambda^{1/2}} \frac{\xi(\nu; A_{+,n}, A_-) d\nu}{(\lambda - \nu^2)^{1/2}} \text{ for a.e. } \lambda > 0, \quad n \in \mathbb{N}. \quad (6.23)$$

A careful investigation in [6] establishes the analog of (6.23) in the limit $n \rightarrow \infty$. However, we emphasize the following formula is not derived in [6] by attempting to take the limit $n \rightarrow \infty$ of either side in (6.23); instead it is derived via a careful application of various trace formulas and the Stieltjes inversion formula resulting in

$$\xi(\lambda; \mathbf{H}_2, \mathbf{H}_1) = \xi(\nu; A_+, A_-) = \frac{1}{2\pi} \int_{\mathbb{R}} dx \phi(x) \quad (6.24)$$

for a.e. $\lambda > 0$ and a.e. $\nu \in \mathbb{R}$. As a consequence of (6.24), the Witten index $W_r(\mathbf{D}_{\mathbf{A}})$ of the non-Fredholm operator $\mathbf{D}_{\mathbf{A}}$ exists and equals

$$W_r(\mathbf{D}_{\mathbf{A}}) = \xi(0_+; \mathbf{H}_2, \mathbf{H}_1) = \xi(0; A_+, A_-) = \frac{1}{2\pi} \int_{\mathbb{R}} dx \phi(x). \quad (6.25)$$

APPENDIX A. SOME FACTS ON SPECTRAL SHIFT FUNCTIONS, TRACE FORMULAS, AND MODIFIED FREDHOLM DETERMINANTS

We recall a few basic facts on spectral shift functions employed in the bulk of this paper and provide results on trace formulas in terms of modified Fredholm determinants.

Closely following [3, Sects. 2–6] and [28, Ch. 8], we provide a brief discussion of how to restrict the open constant in the definition of the spectral shift function $\xi(\cdot; A, A_0)$ up to an integer for a pair of self-adjoint operators (A, A_0) in \mathcal{H} satisfying for some (and hence for all) $z_0 \in \rho(A) \cap \rho(A_0)$,

$$[(A - z_0 I_{\mathcal{H}})^{-1} - (A_0 - z_0 I_{\mathcal{H}})^{-1}] \in \mathcal{B}_1(\mathcal{H}). \quad (\text{A.1})$$

Motivated by the unitary Cayley transforms of A and A_0 , one introduces the modified perturbation determinant,

$$\begin{aligned} \tilde{D}_{A/A_0}(z; z_0) &= \det_{\mathcal{H}}((A - z I_{\mathcal{H}})(A - \bar{z}_0 I_{\mathcal{H}})^{-1}(A_0 - \bar{z}_0 I_{\mathcal{H}})(A_0 - z I_{\mathcal{H}})^{-1}), \\ &z \in \rho(A) \cap \rho(A_0), \operatorname{Im}(z_0) > 0, \end{aligned} \quad (\text{A.2})$$

and notes that (cf. [28, p.270])

$$\overline{\tilde{D}_{A/A_0}(z; z_0)} = \tilde{D}_{A/A_0}(\bar{z}; z_0) / \tilde{D}_{A/A_0}(z_0; z_0), \quad \tilde{D}_{A/A_0}(\bar{z}_0; z_0) = 1, \quad (\text{A.3})$$

and

$$\begin{aligned} \operatorname{tr}_{\mathcal{H}}[(A - z I_{\mathcal{H}})^{-1} - (A_0 - z I_{\mathcal{H}})^{-1}] &= -\frac{d}{dz} \ln(\tilde{D}_{A/A_0}(z; z_0)), \\ &z \in \rho(A) \cap \rho(A_0), \operatorname{Im}(z_0) > 0. \end{aligned} \quad (\text{A.4})$$

In addition,

$$\frac{\tilde{D}_{A/A_0}(z; z_0)}{\tilde{D}_{A/A_0}(\bar{z}; z_0)} = \frac{\tilde{D}_{A/A_0}(z; z_1)}{\tilde{D}_{A/A_0}(\bar{z}; z_1)}, \quad z \in \rho(A) \cap \rho(A_0), \operatorname{Im}(z_0) > 0, \operatorname{Im}(z_1) > 0. \quad (\text{A.5})$$

Then, defining

$$\begin{aligned} \xi(\lambda; A, A_0; z_0) &= (2\pi)^{-1} \lim_{\varepsilon \downarrow 0} [\operatorname{Im}(\ln(\tilde{D}_{A/A_0}(\lambda + i\varepsilon; z_0))) \\ &\quad - \operatorname{Im}(\ln(\tilde{D}_{A/A_0}(\lambda - i\varepsilon; z_0)))] \quad \text{for a.e. } \lambda \in \mathbb{R}, \end{aligned} \quad (\text{A.6})$$

one obtains for $z \in \rho(A) \cap \rho(A_0)$, $\operatorname{Im}(z_0) > 0$, $\operatorname{Im}(z_1) > 0$,

$$\xi(\cdot; A, A_0; z_0) \in L^1(\mathbb{R}; (\lambda^2 + 1)^{-1} d\lambda), \quad (\text{A.7})$$

$$\ln(\tilde{D}_{A/A_0}(z; z_0)) = \int_{\mathbb{R}} \xi(\lambda; A, A_0; z_0) d\lambda [(\lambda - z)^{-1} - (\lambda - \bar{z}_0)^{-1}], \quad (\text{A.8})$$

$$\xi(\lambda; A, A_0; z_0) = \xi(\lambda; A, A_0; z_1) + n(z_0, z_1) \quad \text{for some } n(z_0, z_1) \in \mathbb{Z}, \quad (\text{A.9})$$

$$\operatorname{tr}_{\mathcal{H}}[(A - z)^{-1} - (A_0 - z I_{\mathcal{H}})^{-1}] = - \int_{\mathbb{R}} \frac{\xi(\lambda; A, A_0; z_0) d\lambda}{(\lambda - z)^2}, \quad (\text{A.10})$$

$$[f(A) - f(A_0)] \in \mathcal{B}_1(\mathcal{H}), \quad f \in C_0^\infty(\mathbb{R}), \quad (\text{A.11})$$

$$\operatorname{tr}_{\mathcal{H}}(f(A) - f(A_0)) = \int_{\mathbb{R}} \xi(\lambda; A, A_0; z_0) d\lambda f'(\lambda), \quad f \in C_0^\infty(\mathbb{R}) \quad (\text{A.12})$$

(the final two assertions can be greatly improved).

Up to this point $\xi(\cdot; A, A_0; z_0)$ has been introduced via (A.6) and hence by (A.9), it is determined only up to an additive integer. It is possible to remove this integer

ambiguity in $\xi(\cdot; A, A_0; z_0)$ by adhering to a specific normalization as follows: One introduces

$$U_0(z_0) = (A_0 - z_0 I_{\mathcal{H}})(A_0 - \overline{z_0} I_{\mathcal{H}})^{-1}, \quad U(z_0) = (A - z_0 I_{\mathcal{H}})(A - \overline{z_0} I_{\mathcal{H}})^{-1}, \quad z_0 \in \mathbb{C}_+, \quad (\text{A.13})$$

and then determines a normalized spectral shift function, denoted by $\widehat{\xi}(\cdot; A, A_0; z_0)$, with the help of fixing the branch of $\ln(\widetilde{D}_{A/A_0}(z_0; z_0))$ by the equation,

$$\begin{aligned} i\text{Im}(\ln(\widetilde{D}_{A/A_0}(z_0; z_0))) &= 2i\text{Im}(z_0) \int_{\mathbb{R}} \frac{\widehat{\xi}(\lambda; A, A_0; z_0) d\lambda}{|\lambda - z_0|^2} \\ &= \text{tr}_{\mathcal{H}}(\ln(U(z_0)U_0(z_0)^{-1})). \end{aligned} \quad (\text{A.14})$$

Here $\ln(W)$, with W unitary in \mathcal{H} , is defined via the spectral theorem,

$$W = \oint_{S^1} \mu dE_W(\mu), \quad \ln(W) = i \arg(W) = i \oint_{S^1} \arg(\mu) dE_W(\mu), \quad (\text{A.15})$$

$$\arg(\mu) \in (-\pi, \pi],$$

with $E_W(\cdot)$ the spectral family for W , and

$$\ln(U(z_0)U_0(z_0)^{-1}) = \ln(I_{\mathcal{H}} + [U(z_0) - U_0(z_0)]U_0(z_0)^{-1}) \in \mathcal{B}_1(\mathcal{H}), \quad (\text{A.16})$$

since

$$U(z_0) - U_0(z_0) = -2i\text{Im}(z_0)[(A - \overline{z_0} I_{\mathcal{H}})^{-1} - (A_0 - \overline{z_0} I_{\mathcal{H}})^{-1}] \in \mathcal{B}_1(\mathcal{H}). \quad (\text{A.17})$$

In conjunction with (A.14) we also mention the estimate,

$$\int_{\mathbb{R}} \frac{|\widehat{\xi}(\lambda; A, A_0; z_0)| d\lambda}{|\lambda - z_0|^2} \leq \frac{\pi}{2} \|(A - z_0 I_{\mathcal{H}})^{-1} - (A_0 - z_0 I_{\mathcal{H}})^{-1}\|_{\mathcal{B}_1(\mathcal{H})}. \quad (\text{A.18})$$

Moreover, if there exists $\alpha_0 \geq 0$ such that

$$\alpha \|(A - i\alpha I_{\mathcal{H}})^{-1} - (A_0 - i\alpha I_{\mathcal{H}})^{-1}\|_{\mathcal{B}(\mathcal{H})} < 1, \quad \alpha > \alpha_0, \quad (\text{A.19})$$

then (cf. [28, p. 300–303]),

$$\widehat{\xi}(\cdot; A, A_0; i\alpha) = \widehat{\xi}(\cdot; A, A_0) \text{ is independent of } \alpha \text{ for } \alpha > \alpha_0. \quad (\text{A.20})$$

We also note that if

$$\text{dom}(A) = \text{dom}(A_0), \quad (A - A_0)(A_0 - z_0 I_{\mathcal{H}})^{-1} \in \mathcal{B}_1(\mathcal{H}) \quad (\text{A.21})$$

holds for some (and hence for all) $z_0 \in \rho(A_0)$, then (A.19) is valid for $0 < \alpha_0$ sufficiently large, and (cf. [28, p. 303–304]),

$$\widehat{\xi}(\cdot; A, A_0) \in L^1(\mathbb{R}; (|\lambda| + 1)^{-1-\varepsilon} d\lambda), \quad \varepsilon > 0. \quad (\text{A.22})$$

Since different spectral shift functions only differ by a constant, the inclusion (A.22) remains valid for all spectral shift functions under the assumptions (A.21).

Finally, if

$$B = B^* \in \mathcal{B}_1(\mathcal{H}) \text{ and } A = A_0 + B, \quad (\text{A.23})$$

then (A.19) holds with $\alpha_0 = 0$ and (cf. [28, p. 303–304])

$$\widehat{\xi}(\cdot; A, A_0) \in L^1(\mathbb{R}; d\lambda). \quad (\text{A.24})$$

Assuming (A.23), one usually introduces the (standard) perturbation determinant,

$$\begin{aligned} D_{A/A_0}(z) &= \det_{\mathcal{H}}((A - z I_{\mathcal{H}})(A_0 - z I_{\mathcal{H}})^{-1}) \\ &= \det_{\mathcal{H}}(I_{\mathcal{H}} + B(A_0 - z I_{\mathcal{H}})^{-1}), \quad z \in \rho(A) \cap \rho(A_0), \end{aligned} \quad (\text{A.25})$$

and the associated spectral shift function

$$\xi(\lambda; A, A_0) = \pi^{-1} \lim_{\varepsilon \downarrow 0} \operatorname{Im}(\ln(D_{A/A_0}(\lambda + i\varepsilon))) \text{ for a.e. } \lambda \in \mathbb{R}, \quad (\text{A.26})$$

and hence obtains the following well-known facts for $z \in \rho(A) \cap \rho(A_0)$, $\operatorname{Im}(z_0) > 0$,

$$\overline{D_{A/A_0}(z)} = D_{A/A_0}(\overline{z}), \quad \lim_{|\operatorname{Im}(z)| \rightarrow \infty} D_{A/A_0}(z) = 1, \quad (\text{A.27})$$

$$\xi(\lambda; A, A_0) = (2\pi)^{-1} \lim_{\varepsilon \downarrow 0} [\operatorname{Im}(\ln(D_{A/A_0}(\lambda + i\varepsilon))) - \operatorname{Im}(\ln(D_{A/A_0}(\lambda - i\varepsilon)))]$$

for a.e. $\lambda \in \mathbb{R}$, (A.28)

$$\tilde{D}_{A/A_0}(z; z_0) = D_{A/A_0}(z) / D_{A/A_0}(\overline{z_0}), \quad (\text{A.29})$$

$$\xi(\cdot; A, A_0) \in L^1(\mathbb{R}; d\lambda), \quad (\text{A.30})$$

$$\ln(D_{A/A_0}(z)) = \int_{\mathbb{R}} \xi(\lambda; A, A_0) d\lambda (\lambda - z)^{-1}, \quad (\text{A.31})$$

$$\int_{\mathbb{R}} \xi(\lambda; A, A_0) d\lambda = \operatorname{tr}_{\mathcal{H}}(B), \quad \int_{\mathbb{R}} |\xi(\lambda; A, A_0)| d\lambda \leq \|B\|_{\mathcal{B}_1(\mathcal{H})}, \quad (\text{A.32})$$

Combining the facts (A.6), (A.28), and (A.29) at first only yields for some $n(z_0) \in \mathbb{Z}$,

$$\xi(\lambda; A, A_0; z_0) = \xi(\lambda; A, A_0) + n(z_0) \text{ for a.e. } \lambda \in \mathbb{R}. \quad (\text{A.33})$$

However, also taking into account (A.20) and (A.24) finally yields

$$\widehat{\xi}(\lambda; A, A_0) = \xi(\lambda; A, A_0) \text{ for a.e. } \lambda \in \mathbb{R}. \quad (\text{A.34})$$

Thus, the normalization employed in (A.14) is consistent with the normalization implied by (A.24) in the case of trace class perturbations.

We continue this appendix with the following result, originally derived in [17] under slightly different hypotheses:

Theorem A.1. (i) Suppose A_0 and A are self-adjoint operators with $\operatorname{dom}(A) = \operatorname{dom}(A_0) \subseteq \mathcal{H}$, with $B = \overline{(A - A_0)} \in \mathcal{B}(\mathcal{H})$.

(ii) Assume that for some (and hence for all) $z_0 \in \rho(A_0)$,

$$|B|^{1/2}(A_0 - z_0 I_{\mathcal{H}})^{-1} \in \mathcal{B}_2(\mathcal{H}), \quad (\text{A.35})$$

and that

$$\lim_{z \rightarrow \pm i\infty} \| |B|^{1/2}(A_0 - z I_{\mathcal{H}})^{-1} |B|^{1/2} \|_{\mathcal{B}_2(\mathcal{H})} = 0. \quad (\text{A.36})$$

(iii) Suppose that

$$\operatorname{tr}_{\mathcal{H}}((A_0 - z I_{\mathcal{H}})^{-1} B (A_0 - z I_{\mathcal{H}})^{-1}) = \eta'(z), \quad z \in \rho(A_0), \quad (\text{A.37})$$

where $\eta(\cdot)$ has normal limits, $\lim_{\varepsilon \downarrow 0} \eta(\lambda + i\varepsilon) := \eta(\lambda + i0)$ for a.e. $\lambda \in \mathbb{R}$.

Then

$$\begin{aligned} \int_{\mathbb{R}} \xi(\lambda; A, A_0) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}] &= \eta(z) - \eta(z_0) \\ &+ \ln \left(\frac{\det_{2,\mathcal{H}}(I_{\mathcal{H}} + \operatorname{sgn}(B)|B|^{1/2}(A_0 - z I_{\mathcal{H}})^{-1}|B|^{1/2})}{\det_{2,\mathcal{H}}(I_{\mathcal{H}} + \operatorname{sgn}(B)|B|^{1/2}(A_0 - z_0 I_{\mathcal{H}})^{-1}|B|^{1/2})} \right), \quad z, z_0 \in \rho(A) \cap \rho(A_0), \end{aligned} \quad (\text{A.38})$$

and for some constant $c \in \mathbb{R}$,

$$\begin{aligned} \xi(\lambda; A, A_0) &= \pi^{-1} \operatorname{Im}(\ln(\det_{2,\mathcal{H}}(I_{\mathcal{H}} + \operatorname{sgn}(B)|B|^{1/2}(A_0 - (\lambda + i0)I_{\mathcal{H}})^{-1}|B|^{1/2}))) \\ &\quad + \pi^{-1} \operatorname{Im}(\eta(\lambda + i0)) + c \text{ for a.e. } \lambda \in \mathbb{R}. \end{aligned} \quad (\text{A.39})$$

Remark A.2. (i) We note that Theorem A.1 was derived in [17] for unbounded quadratic form perturbations B of A_0 and hence A_0 was assumed to be bounded from below. Since we here assume that B is bounded, boundedness from below of A_0 is no longer needed and the proof of [17, Theorem 2.3] applies line by line to the current setting. We also note that since $B|_{\operatorname{dom}(A_0)}$ is symmetric and bounded, $B = \overline{(A - A_0)}$ is self-adjoint on \mathcal{H} . The basic identity underlining Theorem A.1 is, of course,

$$\begin{aligned} \operatorname{tr}_{\mathcal{H}}((A - zI_{\mathcal{H}})^{-1} - (A_0 - zI_{\mathcal{H}})^{-1}) + \operatorname{tr}_{\mathcal{H}}((A_0 - zI_{\mathcal{H}})^{-1}B(A_0 - zI_{\mathcal{H}})^{-1}) \quad (\text{A.40}) \\ = -\frac{d}{dz} \ln(\det_{2,\mathcal{H}}(I_{\mathcal{H}} + \operatorname{sgn}(B)|B|^{1/2}(A_0 - zI_{\mathcal{H}})^{-1}|B|^{1/2})), \quad z \in \rho(A) \cap \rho(A_0). \end{aligned}$$

(ii) One can show (cf. (A.51)) that

$$\eta(z) - \eta(z_0) = (z - z_0) \operatorname{tr}_{\mathcal{H}}((A_0 - zI_{\mathcal{H}})^{-1}B(A_0 - z_0I_{\mathcal{H}})^{-1}), \quad z, z_0 \in \rho(A) \cap \rho(A_0). \quad (\text{A.41})$$

We will use this fact later in this appendix. \diamond

For modified Fredholm determinants and their properties, we refer, for instance, to [13, Sect. XI.9], [20, Sect. IV.2] and [25, Ch. 9]. Here we just note that

$$\det_{2,\mathcal{H}}(I_{\mathcal{K}} - A) = \prod_{n \in \mathcal{J}} (1 - \lambda_n(A)) e^{\lambda_n(A)}, \quad A \in \mathcal{B}_2(\mathcal{H}), \quad (\text{A.42})$$

where $\{\lambda_n(A)\}_{n \in \mathcal{J}}$ is an enumeration of the non-zero eigenvalues of A , listed in non-increasing order according to their moduli, and $\mathcal{J} \subseteq \mathbb{N}$ is an appropriate indexing set, and

$$\det_{2,\mathcal{H}}(I_{\mathcal{K}} - A) = \det_{\mathcal{H}}((I_{\mathcal{H}} - A) \exp(A)), \quad A \in \mathcal{B}_2(\mathcal{H}), \quad (\text{A.43})$$

$$\begin{aligned} \det_{2,\mathcal{H}}((I_{\mathcal{K}} - A)(I_{\mathcal{H}} - B)) &= \det_{2,\mathcal{H}}(I_{\mathcal{K}} - A) \det_{2,\mathcal{H}}(I_{\mathcal{H}} - B) e^{-\operatorname{tr}_{\mathcal{H}}(AB)}, \quad (\text{A.44}) \\ A, B &\in \mathcal{B}_2(\mathcal{H}). \end{aligned}$$

In addition, we recall the fact, that $\det_{2,\mathcal{H}}(I_{\mathcal{H}} + \cdot)$ is continuous on $\mathcal{B}_2(\mathcal{H})$, explicitly, for some $c > 0$, one has the estimate (cf. [25, Theorem 9.2(c)])

$$\begin{aligned} |\det_{2,\mathcal{H}}(I_{\mathcal{H}} + T) - \det_{2,\mathcal{H}}(I_{\mathcal{H}} + S)| &\leq \|T - S\|_{\mathcal{B}_2(\mathcal{H})} e^{c[\|S\|_{\mathcal{B}_2(\mathcal{H})} + \|T\|_{\mathcal{B}_2(\mathcal{H})} + 1]^2}, \\ S, T &\in \mathcal{B}_2(\mathcal{H}). \end{aligned} \quad (\text{A.45})$$

In addition, we need some results concerning the connection between trace formulas and modified Fredholm determinants (cf., eg., [20, Sect. IV.2], [25, Ch. 9], [28, Sect. 1.7]). Suppose that A_0 is self-adjoint in the complex, separable Hilbert space \mathcal{H} , and assume that the self-adjoint operator B in \mathcal{H} , satisfies

$$\begin{aligned} \operatorname{dom}(B) &\supseteq \operatorname{dom}(A_0), \\ B &\text{ is infinitesimally bounded with respect to } A_0, \end{aligned} \quad (\text{A.46})$$

and for some (and hence for all) $z_0, z_1 \in \rho(A_0)$,

$$B(A_0 - z_0I_{\mathcal{H}})^{-1} \in \mathcal{B}_2(\mathcal{H}), \quad (A_0 - z_0I_{\mathcal{H}})^{-1}B(A_0 - z_1I_{\mathcal{H}})^{-1} \in \mathcal{B}_1(\mathcal{H}). \quad (\text{A.47})$$

(We recall that B , with $\text{dom}(B) \supseteq \text{dom}(A_0)$, is called infinitesimally bounded with respect to A_0 , if for all $\varepsilon > 0$, there exists $\eta(\varepsilon) > 0$, such that for $f \in \text{dom}(A_0)$, $\|Bf\|_{\mathcal{H}} \leq \varepsilon \|A_0 f\|_{\mathcal{H}} + \eta(\varepsilon) \|f\|_{\mathcal{H}}$.)

Then by assumption (A.46),

$$A = A_0 + B, \quad \text{dom}(A) = \text{dom}(A_0), \quad (\text{A.48})$$

is self-adjoint in \mathcal{H} , and

$$\begin{aligned} [(A - zI_{\mathcal{H}})^{-1} - (A_0 - zI_{\mathcal{H}})^{-1}] &= [(A_0 - zI_{\mathcal{H}})^{-1} B (A_0 - zI_{\mathcal{H}})^{-1}] \\ &\times [(A_0 - zI_{\mathcal{H}})(A - zI_{\mathcal{H}})^{-1}] \in \mathcal{B}_1(\mathcal{H}), \quad z \in \rho(A) \cap \rho(A_0). \end{aligned} \quad (\text{A.49})$$

Given this setup, one concludes the trace formula (cf., e.g., [28, p. 44])

$$\begin{aligned} &\text{tr}_{\mathcal{H}} ((A - zI_{\mathcal{H}})^{-1} - (A_0 - zI_{\mathcal{H}})^{-1}) \\ &= -\frac{d}{dz} \ln(\det_{2,\mathcal{H}}(I_{\mathcal{H}} + B(A_0 - z_0 I_{\mathcal{H}})^{-1})) \\ &\quad - \text{tr}_{\mathcal{H}} ((A_0 - zI_{\mathcal{H}})^{-1} B (A_0 - zI_{\mathcal{H}})^{-1}) \\ &= -\int_{\mathbb{R}} \frac{\xi(\lambda; A, A_0) d\lambda}{(\lambda - z)^2}, \quad z \in \rho(A) \cap \rho(A_0), \end{aligned} \quad (\text{A.50})$$

and consequently, also

$$\begin{aligned} &\ln \left(\frac{\det_{2,\mathcal{H}}(I_{\mathcal{H}} + B(A_0 - zI_{\mathcal{H}})^{-1})}{\det_{2,\mathcal{H}}(I_{\mathcal{H}} + B(A_0 - z_0 I_{\mathcal{H}})^{-1})} \right) \\ &\quad + (z - z_0) \text{tr}_{\mathcal{H}} ((A_0 - zI_{\mathcal{H}})^{-1} B (A_0 - z_0 I_{\mathcal{H}})^{-1}) \\ &= \int_{\mathbb{R}} \xi(\lambda; A, A_0) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}], \quad z, z_0 \in \rho(A) \cap \rho(A_0). \end{aligned} \quad (\text{A.51})$$

(To verify (A.51) it suffices to differentiate either side of (A.51) w.r.t. z , comparing with the final three lines of relation (A.50), and observing that either side of (A.51) vanishes at $z = z_0$.)

At this point we recall that $\det_{2,\mathcal{H}}(I_{\mathcal{H}} + \cdot)$ is continuous on $\mathcal{B}_2(\mathcal{H})$, as recorded earlier in (A.45).

Next, suppose that $A_{0,n}$, B_n , $A_n = A_{0,n} + B_n$, $n \in \mathbb{N}$, and A_0 , B , $A = A_0 + B$ satisfy hypotheses (A.46) and (A.47). Moreover, assume that for some (and hence for all) $z_0 \in \mathbb{C} \setminus \mathbb{R}$,

$$\text{w-lim}_{n \rightarrow \infty} (A_{0,n} - z_0 I_{\mathcal{H}})^{-1} = (A_0 - z_0 I_{\mathcal{H}})^{-1}, \quad (\text{A.52})$$

$$\lim_{n \rightarrow \infty} \|B_n(A_{0,n} - z_0 I_{\mathcal{H}})^{-1} - B(A_0 - z_0 I_{\mathcal{H}})^{-1}\|_{\mathcal{B}_2(\mathcal{H})} = 0, \quad (\text{A.53})$$

$$\begin{aligned} &\lim_{n \rightarrow \infty} \|(A_{0,n} - z_0 I_{\mathcal{H}})^{-1} B_n (A_{0,n} - z_0 I_{\mathcal{H}})^{-1} \\ &\quad - (A_0 - z_0 I_{\mathcal{H}})^{-1} B (A_0 - z_0 I_{\mathcal{H}})^{-1}\|_{\mathcal{B}_1(\mathcal{H})} = 0. \end{aligned} \quad (\text{A.54})$$

One notes that due to self-adjointness of $A_{0,n}$, A_0 , $n \in \mathbb{N}$, relation (A.52) is actually equivalent to strong resolvent convergence, that is,

$$\text{s-lim}_{n \rightarrow \infty} (A_{0,n} - zI_{\mathcal{H}})^{-1} = (A_0 - zI_{\mathcal{H}})^{-1}, \quad z \in \mathbb{C} \setminus \mathbb{R}. \quad (\text{A.55})$$

Moreover, the well-known identity (see, e.g., [26, p. 178]),

$$\begin{aligned} (T_1 - zI_{\mathcal{H}})^{-1} - (T_2 - zI_{\mathcal{H}})^{-1} &= (T_1 - z_0 I_{\mathcal{H}})(T_1 - zI_{\mathcal{H}})^{-1} \\ &\times [(T_1 - z_0 I_{\mathcal{H}})^{-1} - (T_2 - z_0 I_{\mathcal{H}})^{-1}](T_2 - z_0 I_{\mathcal{H}})(T_2 - zI_{\mathcal{H}})^{-1}, \end{aligned} \quad (\text{A.56})$$

$$z, z_0 \in \rho(T_1) \cap \rho(T_2),$$

where T_j , $j = 1, 2$, are linear operators in \mathcal{H} with $\rho(T_1) \cap \rho(T_2) \neq \emptyset$, together with Lemma 3.6, (A.54), and (A.55) imply

$$\begin{aligned} \lim_{n \rightarrow \infty} \left\| (A_{0,n} - z_0 I_{\mathcal{H}})^{-1} B_n (A_{0,n} - z_1 I_{\mathcal{H}})^{-1} \right. \\ \left. - (A_0 - z_0 I_{\mathcal{H}})^{-1} B (A_0 - z_1 I_{\mathcal{H}})^{-1} \right\|_{\mathcal{B}_1(\mathcal{H})} = 0, \quad z_0, z_1 \in \mathbb{C} \setminus \mathbb{R}. \end{aligned} \quad (\text{A.57})$$

Then (A.51) applied to the self-adjoint pairs $(A_n, A_{0,n})$, $n \in \mathbb{N}$, and (A, A_0) , in combination with (A.45)–(A.57) implies the continuity result,

$$\begin{aligned} \lim_{n \rightarrow \infty} \int_{\mathbb{R}} \xi(\lambda; A_n, A_{0,n}) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}] \\ = \lim_{n \rightarrow \infty} \left\{ \ln \left(\frac{\det_{2,\mathcal{H}}(I_{\mathcal{H}} + B_n (A_{0,n} - z I_{\mathcal{H}})^{-1})}{\det_{2,\mathcal{H}}(I_{\mathcal{H}} + B_n (A_{0,n} - z_0 I_{\mathcal{H}})^{-1})} \right) \right. \\ \left. + (z - z_0) \text{tr}_{\mathcal{H}}((A_{0,n} - z I_{\mathcal{H}})^{-1} B_n (A_{0,n} - z_0 I_{\mathcal{H}})^{-1}) \right\} \\ = \ln \left(\frac{\det_{2,\mathcal{H}}(I_{\mathcal{H}} + B (A_0 - z I_{\mathcal{H}})^{-1})}{\det_{2,\mathcal{H}}(I_{\mathcal{H}} + B (A_0 - z_0 I_{\mathcal{H}})^{-1})} \right) \\ + (z - z_0) \text{tr}_{\mathcal{H}}((A_0 - z I_{\mathcal{H}})^{-1} B (A_0 - z_0 I_{\mathcal{H}})^{-1}) \\ = \int_{\mathbb{R}} \xi(\lambda; A, A_0) d\lambda [(\lambda - z)^{-1} - (\lambda - z_0)^{-1}], \quad z, z_0 \in \mathbb{C} \setminus \mathbb{R}. \end{aligned} \quad (\text{A.58})$$

We note that these considerations naturally extend to more complex situations where $A = A_0 +_q B$, $A_n = A_{0,n} +_q B_n$, $n \in \mathbb{N}$, are defined as quadratic form sums of A_0 and B and $A_{0,n}$ and B_n (without assuming any correlation between the domains of A , A_n and A_0), and the modified Fredholm determinants are replaced by symmetrized ones as in Theorem A.1, see, for instance, [16], [17], and [19].

Acknowledgments. We are indebted to Harald Grosse, Jens Kaa, Yuri Latushkin, Matthias Lesch (in particular for sending us [27]), Konstantin Makarov, Alexander Sakhnovich, and Yuri Tomilov for helpful discussions and correspondence.

A.C., F.G., and F.S. thank the Erwin Schrödinger International Institute for Mathematical Physics (ESI), Vienna, Austria, for funding support for this collaboration in form of a Research-in Teams project, “Scattering Theory and Non-Commutative Analysis” for the duration of June 22 to July 22, 2014. F.G. and G.L. are indebted to Gerald Teschl and the ESI for a kind invitation to visit the University of Vienna, Austria, for a period of two weeks in June/July of 2014. A.C., G.L., and F.S. gratefully acknowledge financial support from the Australian Research Council. A.C. also thanks the Alexander von Humboldt Stiftung and colleagues at the University of Münster.

REFERENCES

- [1] N. Anghel, *The two-dimensional magnetic field problem revisited*, J. Math. Phys. **31**, 2091–2093 (1990).
- [2] N. Aronszajn and W. F. Donoghue, *On exponential representations of analytic functions in the upper half-plane with positive imaginary part*, J. Analyse Math. **5**, 321–388 (1956–57).

- [3] M. Sh. Birman and D. R. Yafaev, *The spectral shift function. The work of M. G. Krein and its further development*, St. Petersburg Math. J. **4**, 833–870 (1993).
- [4] D. Bollé, F. Gesztesy, H. Grosse, W. Schweiger, and B. Simon, *Witten index, axial anomaly, and Krein's spectral shift function in supersymmetric quantum mechanics*, J. Math. Phys. **28**, 1512–1525 (1987).
- [5] C. Callias, *Axial anomalies and index theorems on open spaces*, Commun. Math. Phys. **62**, 213–234 (1978).
- [6] A. Carey, F. Gesztesy, G. Levitina, D. Potapov, F. Sukochev, and D. Zanin, *On index theory for non-Fredholm operators: a $(1 + 1)$ -dimensional example*, preprint, 2014, Math. Nachrichten, to appear.
- [7] A. Carey, F. Gesztesy, G. Levitina, D. Potapov, F. Sukochev, and D. Zanin, *Trace formulas for a $(1 + 1)$ -dimensional model operator*, preprint, 2014.
- [8] A. Carey, F. Gesztesy, D. Potapov, F. Sukochev, and Y. Tomilov, *A Jost–Pais-type reduction of Fredholm determinants and some applications*, Integral Eq. Operator Theory **79**, 389–447 (2014).
- [9] A. Carey, F. Gesztesy, D. Potapov, F. Sukochev, and Y. Tomilov, *On the Witten index in terms of spectral shift functions*, ArXiv1404.0740, J. Analyse Math., to appear.
- [10] A. Carey, H. Grosse, and J. KAAD, *Anomalies of Dirac type operators on Euclidean space*, Commun. Math. Phys. **335**, 445–475 (2015).
- [11] A. Carey and J. KAAD, *Topological invariance of the homological index*, preprint arXiv:1402.0475.
- [12] R. W. Carey and J. D. Pincus *An invariant for certain operator algebras*, Proc. Nat. Acad. Sci USA **71**, 1952–1956 (1974).
- [13] N. Dunford and J. Schwartz, *Linear operators. Part II. Spectral theory. Selfadjoint operators in Hilbert space*, Wiley & Sons, New York, 1988.
- [14] F. Gesztesy, Y. Latushkin, K. A. Makarov, F. Sukochev, and Y. Tomilov, *The index formula and the spectral shift function for relatively trace class perturbations*, Adv. Math. **227**, 319–420 (2011).
- [15] F. Gesztesy, Y. Latushkin, F. Sukochev, and Y. Tomilov, *Some operator bounds employing complex interpolation revisited*, Herrnhut Proceedings, W. Arendt, R. Chill, and Y. Tomilov (eds.), Operator Theory: Advances and Applications, Birkhäuser, Springer, Basel (to appear).
- [16] F. Gesztesy and R. Nichols, *Weak convergence of spectral shift functions for one-dimensional Schrödinger operators*, Math. Nachrichten **285**, 1799–1838 (2012).
- [17] F. Gesztesy and R. Nichols, *An abstract approach to weak convergence of spectral shift functions and applications to multi-dimensional Schrödinger operators*, J. Spectral Theory **2**, 225–266 (2012).
- [18] F. Gesztesy and B. Simon, *Topological invariance of the Witten index*, J. Funct. Anal. **79**, 91–102 (1988).
- [19] F. Gesztesy and M. Zinchenko, *Symmetrized perturbation determinants and applications to boundary data maps and Krein-type resolvent formulas*, Proc. London Math. Soc. (3) **104**, 577–612 (2012).
- [20] I. Gohberg and M. G. Krein, *Introduction to the Theory of Linear Nonselfadjoint Operators*, Translations of Mathematical Monographs, Vol. 18, Amer. Math. Soc., Providence, RI, 1969.
- [21] T. Kato, *Perturbation Theory for Linear Operators*, corr. printing of the 2nd ed., Springer, Berlin, 1980.
- [22] A. Pushnitski, *The spectral flow, the Fredholm index, and the spectral shift function*, in *Spectral Theory of Differential Operators: M. Sh. Birman 80th Anniversary Collection*, T. Suslina and D. Yafaev (eds.), AMS Translations, Ser. 2, Advances in the Mathematical Sciences, Vol. 225, Amer. Math. Soc., Providence, RI, 2008, pp. 141–155.
- [23] M. Reed and B. Simon, *Methods of Modern Mathematical Physics. I: Functional Analysis*, revised and enlarged edition, Academic Press, New York, 1980.
- [24] J. Robbin and D. Salamon, *The spectral flow and the Maslov index*, Bull. London Math. Soc. **27**, 1–33 (1995).
- [25] B. Simon, *Trace Ideals and Their Applications*, 2nd ed., Mathematical Surveys and Monographs, Vol. 120, Amer. Math. Soc., Providence, RI, 2005.
- [26] J. Weidmann, *Linear Operators in Hilbert Spaces*, Graduate Texts in Mathematics, Vol. 68, Springer, New York, 1980.

- [27] K. P. Wojciechowski, *Spectral shift function and eta function. Elementary Example*, unpublished manuscript, March, 2007, kindly communicated to us by M. Lesch.
- [28] D. R. Yafaev, *Mathematical Scattering Theory. General Theory*, Amer. Math. Soc., Providence, RI, 1992.

MATHEMATICAL SCIENCES INSTITUTE, AUSTRALIAN NATIONAL UNIVERSITY, KINGSLEY ST., CANBERRA, ACT 0200, AUSTRALIA AND SCHOOL OF MATHEMATICS AND APPLIED STATISTICS, UNIVERSITY OF WOLLONGONG, NSW, AUSTRALIA, 2522

E-mail address: `acarey@maths.anu.edu.au`

URL: `http://maths.anu.edu.au/~acarey/`

DEPARTMENT OF MATHEMATICS, UNIVERSITY OF MISSOURI, COLUMBIA, MO 65211, USA

E-mail address: `gesztesyf@missouri.edu`

URL: `https://www.math.missouri.edu/people/gesztesy`

SCHOOL OF MATHEMATICS AND STATISTICS, UNSW, KENSINGTON, NSW 2052, AUSTRALIA

E-mail address: `g.levitina@student.unsw.edu.au`

SCHOOL OF MATHEMATICS AND STATISTICS, UNSW, KENSINGTON, NSW 2052, AUSTRALIA

E-mail address: `f.sukochev@unsw.edu.au`